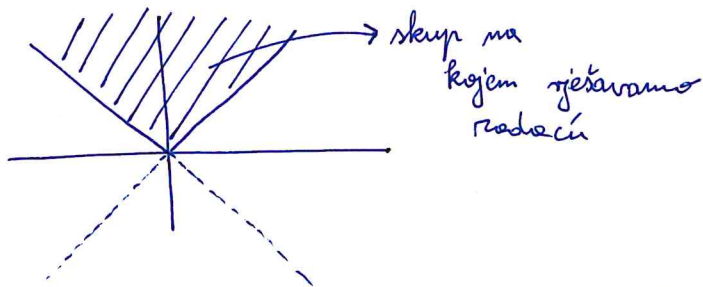


$$2) \begin{cases} \Delta u = 0 \\ u(-x, x) = 0 \\ u(x, x) = x \chi_{[0,1]}(x) \end{cases}$$

$$\text{na } \{(x_1, x_2) \in \mathbb{R}^2 : x_1 + x_2 > 0, x_1 - x_2 < 0\}$$



Pr: 1. NAČIN

Izkoristiti čemu rotacijsku invarijantnost Laplaceovog operatora i svesti problem na domenu $\{(x_1, x_2) \in \mathbb{R}^2 : x_1, x_2 > 0\}$, ~~koji~~ što se već radi na rješavanju.

$$R := \begin{pmatrix} \cos \frac{\pi}{4} & -\sin \frac{\pi}{4} \\ \sin \frac{\pi}{4} & \cos \frac{\pi}{4} \end{pmatrix} = \frac{\sqrt{2}}{2} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} \dots \text{rotacija za } \frac{\pi}{4} \text{ u pozitivnom smjeru}$$

$$v(x_1, x_2) := u(R(x_1, x_2)) = u\left(\frac{\sqrt{2}}{2}(x_1 - x_2, x_1 + x_2)\right)$$

$$\text{tada, } \begin{cases} \Delta v = 0 & \text{u } \{(x_1, x_2) \in \mathbb{R}^2 : x_1, x_2 > 0\} \\ v(0, \cdot) = 0 \\ v(x, 0) = \frac{\sqrt{2}}{2} x \chi_{[0,1]}(\frac{\sqrt{2}}{2} x) = \frac{\sqrt{2}}{2} x \chi_{[0, \sqrt{2}]}(x) \end{cases}$$

Čto je već u neku ruku jednostavniji problem, ali nije identičan zadatku 7 s riješi pa ćemo ipak promatrati.

$$\tilde{v}(x) := v(\sqrt{2}x),$$

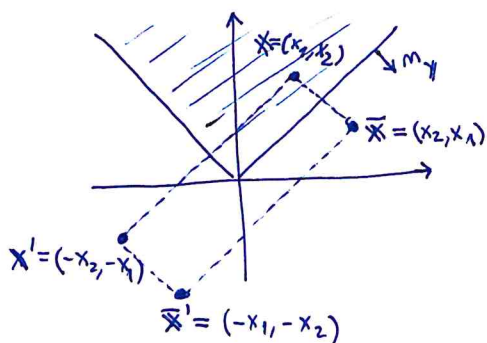
$$\text{pa } \begin{cases} \Delta \tilde{v} = 0 & \text{u } \{(x_1, x_2) \in \mathbb{R}^2 : x_1, x_2 > 0\} \\ \tilde{v}(0, \cdot) = 0 \\ \tilde{v}(x, 0) = x \chi_{[0,1]}(x) \end{cases}$$

tada s riješi računom:

$$\tilde{v}(x_1, x_2) = \frac{x_2}{\pi} \ln \sqrt{\frac{(x_1-1)^2 + x_2^2}{(x_1+1)^2 + x_2^2}} - \frac{x_1}{\pi} \left(\operatorname{arctg} \left(\frac{x_1-1}{x_2} \right) - \operatorname{arctg} \left(\frac{x_1+1}{x_2} \right) \right).$$

$$\begin{aligned} u(x_1, x_2) &= v(R^T(x_1, x_2)) \\ &= \tilde{v}\left(\frac{\sqrt{2}}{2} \cdot \frac{\sqrt{2}}{2} \begin{pmatrix} 1 & 1 \\ -1 & 1 \end{pmatrix} (x_1, x_2)\right) \\ &= \tilde{v}\left(\frac{x_1+x_2}{2}, \frac{x_2-x_1}{2}\right) \\ &= \frac{x_2-x_1}{2\pi} \ln \sqrt{\frac{(x_1+x_2-2)^2 + (x_2-x_1)^2}{(x_1+x_2+2)^2 + (x_2-x_1)^2}} - \frac{x_1+x_2}{2\pi} \left(\operatorname{arctg} \frac{x_1+x_2-2}{x_2-x_1} - \operatorname{arctg} \frac{x_1+x_2+2}{x_2-x_1} \right) \end{aligned}$$

2. NAČIN Obrediti Greenov f-ju konstanti refleksije



$$\Rightarrow G(x, y) = \phi(|x-y|) - \phi(|\bar{x}-y|) - \phi(|x'-y|) + \phi(|\bar{x}'-y|)$$

$$= -\frac{1}{4\pi} \ln((x_1-y_1)^2 + (x_2-y_2)^2) + \frac{1}{4\pi} \ln((x_2-y_1)^2 + (x_1-y_2)^2) + \frac{1}{4\pi} \ln((x_2+y_1)^2 + (x_1+y_2)^2) - \frac{1}{4\pi} \ln((x_1+y_1)^2 + (x_2+y_2)^2)$$

vrednost f-je na robu

$$\Rightarrow u(x_1, x_2) = - \int_{\{(x_1, x_2): x_1=x_2, x_1>0\}} \frac{\partial G}{\partial n_y}(x, y) \overbrace{g(y)}^{vrednost f-je na robu} dS_y$$

$$= \left\{ \begin{array}{l} \gamma: [0, \infty) \rightarrow \{(x_1, x_2): x_1=x_2, x_1>0\} \\ \gamma(y) = (y, y) \\ \gamma'(y) = (1, 1) \Rightarrow |\gamma'(y)| = \sqrt{2} \end{array} \right\}$$

$$= - \int_0^\infty \frac{\partial G}{\partial n_y}(x_1, x_2, \gamma(y)) g(\gamma(y)) |\gamma'(y)| dy$$

$$= -\sqrt{2} \int_0^1 \frac{\partial G}{\partial n_y}(x_1, x_2, y, y) y dy$$

$$n(y, y) = \frac{\sqrt{2}}{2}(1, -1) \dots \text{jedinična normala}$$

$$\Rightarrow \frac{\partial G}{\partial n_y}(x_1, x_2, y, y) = \frac{\sqrt{2}}{2} \frac{\partial G}{\partial y_1}(x, y, y) - \frac{\sqrt{2}}{2} \frac{\partial G}{\partial y_2}(x, y, y)$$

$$\Rightarrow u(x_1, x_2) = \int_0^1 \frac{\partial G}{\partial y_2}(x_1, x_2, y, y) y dy - \int_0^1 \frac{\partial G}{\partial y_1}(x_1, x_2, y, y) y dy$$

Ako je \tilde{G} Greenova f-ja na $\{(x_1, x_2) \in \mathbb{R}^2 : x_1, x_2 > 0\}$ tada se lahko pokaže da je

$$G(x, y) = \tilde{G}(R^T x, R^T y),$$

pri čemer je

$$R = \frac{\sqrt{2}}{2} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix} \quad (\text{rot. na } \frac{\pi}{4})$$

$$3) \begin{cases} u_t - \Delta u = 0 & \text{u } \mathbb{R}^+ \times \mathbb{R}^2 \\ u(0, x_1, x_2) = \sin(2x_1 - x_2) \end{cases}$$

Pj. Po formuli za rješenje imamo

$$\begin{aligned} u(t, x_1, x_2) &= \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} \Phi(t, x_1 - y_1, x_2 - y_2) \overbrace{\sin(2y_1 - y_2)}^{\sin 2y_1 \cos y_2 - \cos 2y_1 \sin y_2} dy_1 dy_2 \\ &= \frac{1}{4\pi t} \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} e^{-\frac{(x_1 - y_1)^2}{4t}} e^{-\frac{(x_2 - y_2)^2}{4t}} (\sin 2y_1 \cos y_2 - \cos 2y_1 \sin y_2) dy_1 dy_2 \\ &= \frac{1}{4\pi t} \int_{-\infty}^{\infty} \cos y_2 e^{-\frac{(x_2 - y_2)^2}{4t}} \left(\int_{-\infty}^{\infty} e^{-\frac{(x_1 - y_1)^2}{4t}} \sin 2y_1 dy_1 \right) dy_2 \\ &\quad - \frac{1}{4\pi t} \int_{-\infty}^{\infty} \sin y_2 e^{-\frac{(x_2 - y_2)^2}{4t}} \left(\int_{-\infty}^{\infty} e^{-\frac{(x_1 - y_1)^2}{4t}} \cos 2y_1 dy_1 \right) dy_2 \end{aligned}$$

$$\begin{cases} z = y_1 - x_1 \Rightarrow dz = dy_1 \\ y_1 = z + x_1 \end{cases}$$

$$\begin{aligned} &= \frac{1}{4\pi t} \int_{-\infty}^{\infty} \cos y_2 e^{-\frac{(x_2 - y_2)^2}{4t}} \left(\int_{-\infty}^{\infty} e^{-\frac{z^2}{4t}} \overbrace{(\sin 2z \cos 2x_1 + \sin 2x_1 \cos 2z)}^{\text{ovaj dio bude 0}} dz \right) dy_2 \\ &\quad - \frac{1}{4\pi t} \int_{-\infty}^{\infty} \sin y_2 e^{-\frac{(x_2 - y_2)^2}{4t}} \left(\int_{-\infty}^{\infty} e^{-\frac{z^2}{4t}} \underbrace{(\cos 2z \cos 2x_1 - \sin 2z \sin 2x_1)}_{\text{ovaj dio bude 0}} dz \right) dy_2 \end{aligned}$$

$$\begin{aligned} &= \frac{1}{4\pi t} \int_{-\infty}^{\infty} \cos y_2 e^{-\frac{(x_2 - y_2)^2}{4t}} \left(\int_{-\infty}^{\infty} e^{-\frac{z^2}{4t}} \cos 2z dz \right) \sin 2x_1 dy_2 \\ &\quad - \frac{1}{4\pi t} \int_{-\infty}^{\infty} \sin y_2 e^{-\frac{(x_2 - y_2)^2}{4t}} \left(\int_{-\infty}^{\infty} e^{-\frac{z^2}{4t}} \cos 2z dz \right) \cos 2x_1 dy_2 \end{aligned}$$

$$\begin{aligned} &= \frac{\sin 2x_1}{\sqrt{4\pi t}} e^{-4t} \int_{-\infty}^{\infty} e^{-\frac{(x_2 - y_2)^2}{4t}} \cos y_2 dy_2 - \frac{\cos 2x_1}{\sqrt{4\pi t}} e^{-4t} \int_{-\infty}^{\infty} e^{-\frac{(x_2 - y_2)^2}{4t}} \sin y_2 dy_2 \end{aligned}$$

$$\begin{cases} z = y_2 - x_2 \Rightarrow dz = dy_2 \\ y_2 = z + x_2 \end{cases}$$

$$\begin{aligned}
&= \frac{\sin 2x_1}{\sqrt{4\pi t}} e^{-4t} \int_{-\infty}^{\infty} e^{-\frac{z^2}{4t}} (\cos z \cos x_2 - \underbrace{\sin z \sin x_2}_{=0}) dz \\
&\quad - \frac{\cos 2x_1}{\sqrt{4\pi t}} e^{-4t} \int_{-\infty}^{\infty} e^{-\frac{z^2}{4t}} (\underbrace{\sin z \cos x_2}_{=0} + \cos z \sin x_2) dz \\
&= \frac{\sin 2x_1 \cos x_2}{\sqrt{4\pi t}} e^{-4t} \int_{-\infty}^{\infty} e^{-\frac{z^2}{4t}} \cos z dz - \frac{\cos 2x_1 \sin x_2}{\sqrt{4\pi t}} e^{-4t} \int_{-\infty}^{\infty} e^{-\frac{z^2}{4t}} \cos z dz \\
&= \sin 2x_1 \cos x_2 e^{-4t} e^{-t} - \cos 2x_1 \sin x_2 e^{-4t} e^{-t} \\
&= e^{-5t} (\sin 2x_1 \cos x_2 - \cos 2x_1 \sin x_2) \\
&= \underline{e^{-5t} \sin(2x_1 - x_2)}
\end{aligned}$$

NAP. Naravno, ovo rješenje se moglo i pogoditi pa bismo bilo potrebno
povijesti da radoroljiva početnu radeću.

$$4) \begin{cases} u_{tt} - \Delta u = t+1 & u \in \mathbb{R}^+ \times \mathbb{R}^2 \\ u(0, x_1, x_2) = 0 \\ u_t(0, x_1, x_2) = x_1^2 \end{cases}$$

R: Rastant čeno problem ne dno:

$$\textcircled{1} \begin{cases} u_{tt} - \Delta u = 0 \\ u(0, x_1, x_2) = 0 \\ u_t(0, x_1, x_2) = x_1^2 \end{cases}, \quad \textcircled{2} \begin{cases} u_{tt} - \Delta u = t+1 \\ u(0, \cdot) = 0 \\ u_t(0, \cdot) = 0 \end{cases}$$

Por Poissonovj formuli na $\textcircled{1}$ imamo:

$$u_1(t, x_1, x_2) = \frac{1}{2\pi} \int_{\mathbb{R}^2} \frac{t^2 y_1^2}{K(x, t) (t^2 - |y - x|^2)^{1/2}} dy$$

$$= \left\{ \begin{array}{l} y_1 = r \cos \varphi + x_1 \\ y_2 = r \sin \varphi + x_2 \end{array} \rightarrow \text{Jacobijan je } r \right\}$$

$$= \frac{1}{2\pi} \int_0^t \int_0^{2\pi} r \frac{(r \cos \varphi + x_1)^2}{(t^2 - r^2)^{1/2}} d\varphi dr$$

$$= \frac{1}{2\pi} \int_0^t \frac{r}{(t^2 - r^2)^{1/2}} \left(\int_0^{2\pi} r^2 \cos^2 \varphi + 2rx_1 \cos \varphi + x_1^2 d\varphi \right) dr$$

$$\left\{ \begin{array}{l} \bullet \int_0^{2\pi} \cos^2 \varphi d\varphi = \int_0^{2\pi} \frac{\cos(2\varphi) + 1}{2} d\varphi \\ \bullet \int_0^{2\pi} \cos \varphi d\varphi = \sin \varphi \Big|_0^{2\pi} = 0 \end{array} \right.$$

$$= \frac{\sin(2\varphi)}{4} \Big|_0^{2\pi} + \frac{1}{2} \cdot 2\pi = \pi$$

$$= \frac{1}{2} \int_0^t \frac{r^3}{(t^2 - r^2)^{1/2}} dr + x_1^2 \int_0^t \frac{r}{(t^2 - r^2)^{1/2}} dr$$

$$= \left\{ \begin{array}{l} s = t^2 - r^2 \Rightarrow r^2 = t^2 - s \\ ds = -2r dr \\ r=0 \Rightarrow s=t^2, r=t \Rightarrow s=0 \end{array} \right.$$

$$= \frac{1}{4} \int_0^{t^2} \frac{(t^2 - s)}{\sqrt{s}} ds + \frac{x_1^2}{2} \int_0^{t^2} \frac{ds}{\sqrt{s}} = \left(\frac{x_1^2}{2} + \frac{t^2}{4} \right) \int_0^{t^2} \frac{ds}{\sqrt{s}} - \frac{1}{4} \int_0^{t^2} s^{1/2} ds$$

$$= \left(x_1^2 + \frac{t^2}{2} \right) \sqrt{s} \Big|_0^{t^2} - \frac{1}{4} \cdot \frac{2}{3} s^{3/2} \Big|_0^{t^2}$$

$$= t x_1^2 + \frac{t^3}{2} - \frac{1}{6} t^3 = t x_1^2 + \frac{t^3}{3}$$

Za ② koristimo D'Alembertovo načelo pa prometramo

$$\begin{cases} v_{tt}(\cdot, \cdot; \Delta) - \Delta v(\cdot, \cdot; \Delta) = 0 \\ v(\Delta, \cdot; \Delta) = 0 \\ v_t(\Delta, \cdot; \Delta) = \Delta + 1 \end{cases}$$

Onda $\tilde{v}(t, \mathbf{x}) := v(t + \Delta, \mathbf{x}; \Delta)$ zadovoljava

$$\begin{cases} \tilde{v}_{tt} - \Delta \tilde{v} = 0 \\ \tilde{v}(0, \cdot) = 0 \\ \tilde{v}_t(0, \cdot) = \Delta + 1 \end{cases}$$

Pa Poissonovoj formuli imamo

$$\tilde{v}(t, x_1, x_2) = \frac{\Delta + 1}{2\pi} \int_{K(\mathbf{x}, t)} \frac{dy}{(t^2 - |\mathbf{y} - \mathbf{x}|^2)^{1/2}} = \begin{cases} y_1 = r \cos \varphi + x_1 \\ y_2 = r \sin \varphi + x_2 \end{cases}$$

$$= \frac{\Delta + 1}{2\pi} \left(\int_0^t \frac{r}{(t^2 - r^2)^{1/2}} dr \right) \left(\int_0^{2\pi} d\varphi \right)$$

$$= (\Delta + 1) \int_0^t \frac{r}{(t^2 - r^2)^{1/2}} dr = t(\Delta + 1)$$

$$\Rightarrow v(t, x_1, x_2; \Delta) = \tilde{v}(t - \Delta, x_1, x_2) = (t - \Delta)(\Delta + 1) = (t - 1)\Delta - \Delta^2 + t$$

$$\Rightarrow u_2(t, x_1, x_2) = \int_0^t v(t, x_1, x_2; \Delta) d\Delta = - \int_0^t \Delta^2 d\Delta + (t - 1) \int_0^t \Delta d\Delta + t \int_0^t d\Delta$$

$$= - \frac{t^3}{3} + (t - 1) \frac{t^2}{2} + t^2$$

$$= - \frac{t^3}{3} + \frac{t^3}{2} - \frac{t^2}{2} + t^2 = \frac{t^3}{6} + \frac{t^2}{2}$$

$$\Rightarrow u(t, x_1, x_2) = t x_1^2 + \frac{t^3}{3} + \frac{t^3}{6} + \frac{t^2}{2}$$

$$= t x_1^2 + \frac{t^3}{2} + \frac{t^2}{2}$$

NAP. Naravno, tj. ①, ②
ili cijelog problema se
moglo i pogoditi.