Homogenization of Random Multilevel Junction

G.A.Chechkin, T.P.Chechkina

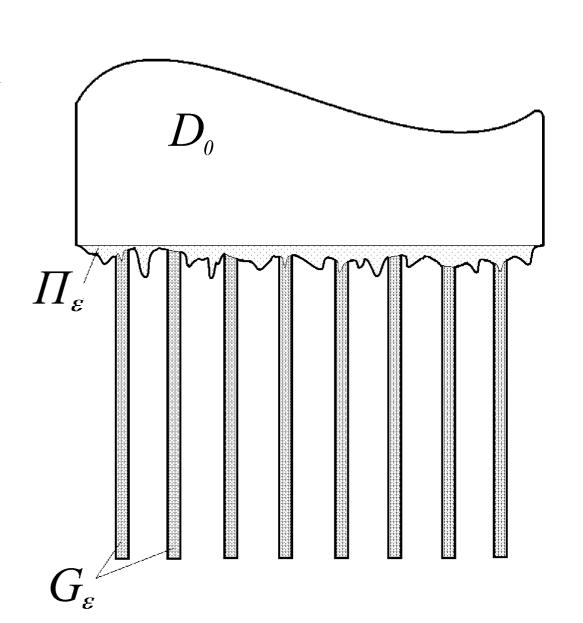
Moscow Lomonosov State University & Narvik University College

and

Moscow Engineering Physical Institute (State University),

This work was done in collaboration with Taras Mel'nyk (Kyev, Ukraine), Ciro D'Apice (Salerno, Italy) and Umberto De Maio (Naples, Italy).

Figure 1: Cascade thick junction with random transmission zone

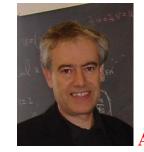


 D_{ε}





NAZAROV S.A.



AMIRAT Y.



BODART O.



DE MAIO U.



GAUDIELLO A.



Geometrical description

Let a, h < 1 be positive real numbers, $I_h(1/2) := \left(\frac{1-h}{2}, \frac{1+h}{2}\right)$ belong to (0, 1). The segment $I_0 := [0, a]$ consists of N subsegments $[\varepsilon j, \varepsilon (j+1)], j = 0, \dots, N-1$. Here $\varepsilon = a/N$ is a small parameter.

Cascade thick junction D_{ε} with random transmission zone consists of a body

$$D_0 = \{ x \in \mathbb{R}^2 : 0 < x_1 < a, \quad 0 < x_2 < \Phi(x_1) \},\$$

 $\Phi \in C^1([0,a]), \quad \min_{[0,a]} \Phi > 0; \text{ thin rectangles}$

$$\widehat{G}_{j}(\varepsilon) = \left\{ x \in \mathbb{R}^{2} : (x_{1}, 0) \in I_{0}, \left| x_{1} - \varepsilon \left(j + \frac{1}{2} \right) \right| < \frac{\varepsilon h}{2}, \ x_{2} \in (-l, 0) \right\}, j = 0, \dots, N - 1$$

and thin layer with oscillating boundary

$$\Pi_{\varepsilon} = \left\{ x \in \mathbb{R}^2 : x_1 \in (0, a), \varepsilon \theta(x_1) F\left(\frac{x_1}{\varepsilon}, \omega\right) < x_2 \le 0 \right\},\$$

Geometrical description

where $\theta(x_1)$ is a smooth nonnegative function with $supp \ \theta(x_1) \subset I_0$ and $F(\xi_1, \omega)$ is a random statistically homogeneous ergodic nonpositive function with smooth realizations, ω is an element of a standard probability space $(\Omega, \mathcal{A}, \mu)$. Thus,

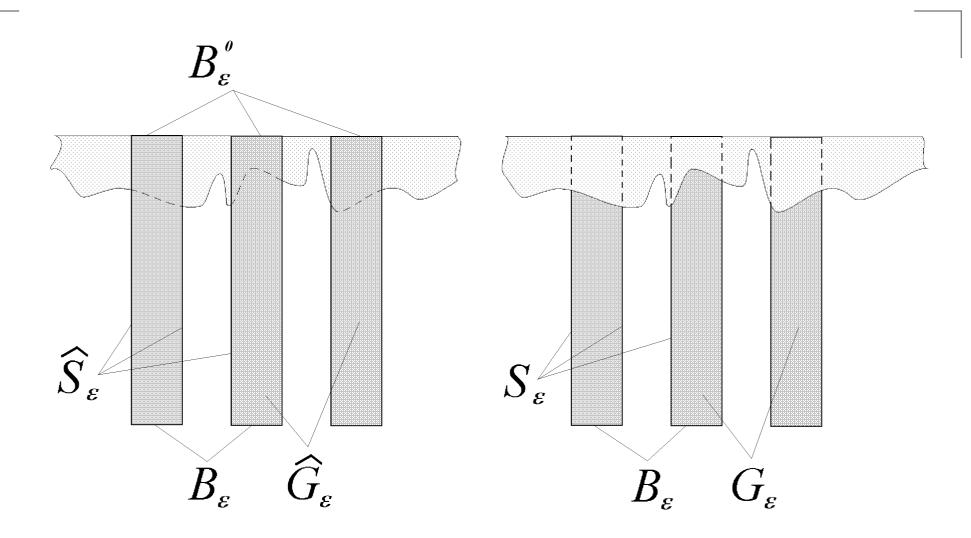
$$D_{\varepsilon} = D_0 \cup \Pi_{\varepsilon} \cup \widehat{G}_{\varepsilon},$$

where

$$\widehat{G}_{\varepsilon} = \bigcup_{j=0}^{N-1} \widehat{G}_j(\varepsilon)$$

or $D_{\varepsilon} = D_0 \cup \Pi_{\varepsilon} \cup G_{\varepsilon}$, where $G_{\varepsilon} = \hat{G}_{\varepsilon} \setminus \Pi_{\varepsilon}$. We denote also $B_{\varepsilon}^0 = \left\{ x \in \mathbb{R}^2 : |x_1 - \varepsilon(j + \frac{1}{2})| < \frac{\varepsilon h}{2}, x_2 = 0 \right\}$, $\Gamma_{\varepsilon} = \left\{ x \in \mathbb{R}^2 : x_1 \in (0, a) \setminus B_{\varepsilon}^0, \ \varepsilon \theta(x_1) F\left(\frac{x_1}{\varepsilon}, \omega\right) = x_2 \right\}$, $\hat{\Upsilon}_{\varepsilon} := \partial \hat{G}_{\varepsilon} \setminus \overline{B}_{\varepsilon}^0$ or $\hat{\Upsilon}_{\varepsilon} = \hat{S}_{\varepsilon} \cup B_{\varepsilon}$, where \hat{S}_{ε} is the lateral surface of the set $\hat{G}_{\varepsilon}, B_{\varepsilon}$ is the lower surface of \hat{G}_{ε} ; $\Upsilon_{\varepsilon} := \partial G_{\varepsilon} \setminus \partial \Pi_{\varepsilon}, \ \Upsilon_{\varepsilon} = S_{\varepsilon} \cup B_{\varepsilon}, \ \Gamma_1 = \{x : x_2 = \Phi(x_1), \ x_1 \in [0, a]\},$ $\gamma_{\varepsilon} = \partial D_{\varepsilon} \setminus (\Gamma_{\varepsilon} \cup \Upsilon_{\varepsilon} \cup \Gamma_1).$

Figure 2: Rectangles and random layer



Setting of the problem

In D_{ε} we consider the following problem:

$$-\Delta_{x} u_{\varepsilon}(x) = f_{\varepsilon}(x), \qquad x \in D_{\varepsilon};$$

$$\partial_{\nu} u_{\varepsilon}(x) + \varepsilon^{\tau} \theta(x_{1}) p\left(\frac{x_{1}}{\varepsilon}, \omega\right) u_{\varepsilon}(x) = \theta(x_{1}) q\left(\frac{x_{1}}{\varepsilon}, \omega\right), \qquad x \in \Gamma_{\varepsilon};$$

$$\partial_{\nu} u_{\varepsilon}(x) + \varepsilon^{\mu} k u_{\varepsilon}(x) = \varepsilon^{\beta} g_{\varepsilon}(x), \qquad x \in \Upsilon_{\varepsilon};$$

$$u_{\varepsilon}(x) = 0, \qquad x \in \Gamma_{1};$$

$$\partial_{\nu} u_{\varepsilon}(x) = 0, \qquad x \in \gamma_{\varepsilon}.$$

(1)

Here $\partial_{\nu} = \partial/\partial \nu$ is the derivative with respect to the outer normal; the constant k is positive; the parameters $\beta \ge 1$, μ , τ are real; $p(\xi_1, \omega)$ and $q(\xi_1, \omega)$ are random statistically homogeneous ergodic positive functions. The functions p and q have smooth realizations. It should be noted that

$$[u_{\varepsilon}] = 0, \qquad [\partial_{x_2} u_{\varepsilon}] = 0 \qquad \text{on} \quad I_0 \cap \Pi_{\varepsilon},$$

where $[\cdot]$ is the jump of a function.

Assume that $(\Omega, \mathcal{A}, \mu)$ is a probability space, i.e. the set Ω with σ -algebra \mathcal{A} of its subsets and σ -additive nonnegative measure μ on \mathcal{A} such that $\mu(\Omega) = 1$.

Definition 1. A family of measurable maps

 $T_{x_1}: \Omega \to \Omega, \quad x_1 \in \mathbb{R}$

we call a dynamical system, if the following properties hold true:

group property:

$$T_{x_1+y_1} = T_{x_1}T_{y_1} \quad \forall x_1, y_1 \in \mathbb{R}; \qquad T_0 = Id$$

(*Id* is the identical mapping);

isometry property (the mapping T_{x_1} preserves the measure μ on Ω):

$$T_{x_1}A \in \mathcal{A}, \quad \mu(T_{x_1}A) = \mu(A) \quad \forall x_1 \in \mathbb{R}, \ A \in \mathcal{A};$$

measurability: for any measurable functions $\Psi(\omega)$ on Ω the function $\Psi(T_{x_1}\omega)$ is measurable on $\Omega \times \mathbb{R}$ and continuous in x_1 .

Let $L^q(\Omega,\mu)$ $(q \ge 1)$ be the space of measurable functions integrable in the power q with respect to the measure μ . If $U_{x_1}: \Omega \to \Omega$ is a dynamical system, then in the space $L^2(\Omega,\mu)$ we define a parametric group of operators $\{U_{x_1}\}, x_1 \in \mathbb{R}$ (we keep the same notation) by the formula

$$(U_{x_1}\Psi)(\omega) := \Psi(U_{x_1}\omega), \quad \Psi \in L^2(\Omega,\mu).$$

From the condition 3) of the definition it follows that the group U_{x_1} is strongly continuous, i.e. for any $\Psi \in L^2(\Omega, \mu)$

$$\lim_{x_1 \to 0} \|U_{x_1} \Psi - \Psi\|_{L^2(\Omega, \mu)} = 0.$$

Definition 2. Suppose that $\Psi(\omega)$ is a measurable function on Ω . The function $\Psi(T_{x_1}\omega)$ of $x_1 \in \mathbb{R}$ for fixed $\omega \in \Omega$ is called the realization of the function Ψ .

Proposition 1. Assume that $\Psi \in L^q(\Omega, \mu)$, then almost all realizations $\Psi(T_{x_1}\omega)$ belong to $L^q_{loc}(\mathbb{R})$. If the sequence $\Psi_k \in L^q(\Omega, \mu)$ converges in $L^q(\Omega, \mu)$ to the function Ψ , then there exists a subsequence k' such that almost all realizations $\Psi_{k'}(T_{x_1}\omega)$ converges in $L^q_{loc}(\mathbb{R})$ to the realization $\Psi(T_{x_1}\omega)$. **Definition 3.** A measurable function $\Psi(\omega)$ on Ω is called invariant, if $\Psi(T_{x_1}\omega) = \Psi(\omega)$ for any $x_1 \in \mathbb{R}$ and almost all $\omega \in \Omega$.

Definition 4. The dynamical system T_{x_1} is called ergodic, if any invariant function almost everywhere coincides with a constant.

We denote by \mathcal{B} the natural Borel σ -algebra of subsets of the space \mathbb{R} . Suppose that $F(x_1) \in L^1_{loc}(\mathbb{R})$.

Definition 5. We say that the function $\mathbf{F}(x_1)$ has a spatial average, if the limit

$$M(F) = \lim_{\varepsilon \to 0} \frac{1}{|B|} \int_{B} F\left(\frac{x_1}{\varepsilon}\right) \, dx_1 \tag{2}$$

does exist for any bounded Borel sets $B \in \mathcal{B}$ and does not depend on the choice of B, and M(F) is called the spatial average value of the function F.

In equivalent form

$$M(\mathbf{F}) = \lim_{t \to +\infty} \frac{1}{|B_t|} \int_{B_t} \mathbf{F}(x_1) \, dx_1, \tag{3}$$

where $B_t = \{x \in \mathbb{R} \mid \frac{x_1}{t} \in B\}.$

Proposition 2. Let the function $F(x_1)$ have a spatial meanvalue in \mathbb{R} , and suppose that the family $\{F(\frac{x_1}{\varepsilon}), 0 < \varepsilon \leq 1\}$ is bounded in $L^q(\mathcal{K})$ for some $q \geq 1$, where \mathcal{K} is a compact in \mathbb{R} . Then

$$F\left(\frac{x_1}{\varepsilon}\right) \to M(F)$$
 weakly in $L^q_{loc}(\mathbb{R})$ as $\varepsilon \to 0$.

Proposition 3. (Birkhoff ergodic theorem) Let T_{x_1} satisfy the Definition 1 and assume that $\Psi \in L^q(\Omega, \mu), q \geq 1$. Then for almost all $\omega \in \Omega$ the realization $\Psi(T_{x_1}\omega)$ has the spatial meanvalue $M(\Psi(T_{x_1}\omega))$. Moreover, the spatial meanvalue $M(\Psi(T_{x_1}\omega))$ is a conditional mathematical expectation of the function $\Psi(\omega)$ with respect to the σ -algebra of invariant subsets. Hence, $M(\Psi(T_{x_1}\omega))$ is an invariant function and

$$\mathbb{E}(\Psi) \equiv \int_{\Omega} \Psi(\omega) \, d\mu = \int_{\Omega} M(\Psi(T_{x_1}\omega)) \, d\mu. \tag{4}$$

In particular, if the dynamical system T_{x_1} is ergodic, then for almost all $\omega \in \Omega$ the following formula

$$\mathbb{E}(\Psi) = M(\Psi(T_{x_1}\omega))$$

holds true.

Definition 6. A random function $\Psi(x_1, \omega)$ $(x_1 \in \mathbb{R}, \omega \in \Omega)$ is called statistically homogeneous, if the following representation

$$\Psi(x_1,\omega) = \Psi(T_{x_1}\omega)$$

holds for some function Ψ , where T_{x_1} is a dynamical system in Ω .

For statistically homogeneous functions with smooth realizations we denote

$$\partial_{\omega} \Psi(T_{x_1}\omega) := \partial_{x_1} \Psi(x_1,\omega)$$

Assumptions

We assume that the following conditions are fulfilled. Without loss of generality $f_{\varepsilon} \in L^2(D_1)$, where $\overline{D}_1 = \overline{D}_0 \cup \overline{D}_2$, $D_2 = (0, a) \times (-l, 0)$, and

 $f_{\varepsilon} \to f_0$ strongly in $L^2(D_1)$ as $\varepsilon \to 0.$ (5)

The function $g_{\varepsilon} \in H^1(D_2)$ and

 $g_{\varepsilon} \rightharpoonup g_0$ weakly in $H^1(D_2)$ as $\varepsilon \to 0.$ (6)

Now let us formulate the conditions for functions p, q and F. We assume that p, q and F are statistically homogeneous, T_{x_1} is ergodic and a.s.

 $\mathbf{p}(\omega) \ge 0, \qquad \|\mathbf{p}\|_{L^{\infty}(\Omega,\mu)} < \infty, \qquad \|\mathbf{q}\|_{L^{\infty}(\Omega,\mu)} < \infty,$

 $\|\mathbf{F}\|_{L^{\infty}(\Omega,\mu)} < \infty, \qquad \|\partial_{\omega}\mathbf{F}\|_{L^{\infty}(\Omega,\mu)} < \infty.$

We define the continuation by zero for functions from $H^1(G_{\varepsilon})$ in the following manner:

$$\widetilde{y_{\varepsilon}}(x) = \begin{cases} y_{\varepsilon}, & x \in G_{\varepsilon}, \\ 0, & x \in D_2 \setminus G_{\varepsilon}, \end{cases}$$

where $D_2 = (0, a) \times (-l, 0)$.

Theorem 1 (The case $\tau \ge 0$ and $\mu \ge 1$). The solution u_{ε} to the problem (1) for almost all ω (a.s.) satisfies

$$\begin{split} u_{\varepsilon} & \rightharpoonup v_{0}^{+} & \text{in } H^{1}(D_{0},\Gamma_{1}), \qquad \widetilde{u_{\varepsilon}} & \rightharpoonup h v_{0}^{-} & \text{in } L^{2}(D_{2}), \\ \widetilde{\partial_{x_{2}}u_{\varepsilon}} & \rightharpoonup h \partial_{x_{2}}v_{0}^{-} & \text{in } L^{2}(D_{2}), \qquad \widetilde{\partial_{x_{1}}u_{\varepsilon}} & \rightharpoonup 0 & \text{in } L^{2}(D_{2}), \end{split}$$
(7)

$$as \varepsilon \to 0, \text{ where the function } v_{0}(x) = \begin{cases} v_{0}^{+}(x), & x \in D_{0}, \\ v_{0}^{-}(x), & x \in D_{2}, \end{cases} \text{ is the unique solution to the problem} \\ \begin{cases} -\Delta_{x} v_{0}^{+}(x) = f_{0}(x), & x \in D_{0} \\ v_{0}^{+}(x) = 0, & x \in \Gamma_{1} \\ \partial_{\nu}v_{0}^{+}(x) = 0, & x \in \partial D_{0} \setminus (\Gamma_{1} \cup I_{0}), \\ -h\partial_{x_{2}x_{2}}^{2}v_{0}^{-}(x) + 2\delta_{\mu,1}kv_{0}^{-}(x) = hf_{0}(x) + \delta_{\beta,1}g_{0}(x), \quad x \in D_{2}, \\ v_{0}^{+}(x_{1}, 0) = v_{0}^{-}(x_{1}, 0), & (x_{1}, 0) \in I_{0}, \\ (h\partial_{x_{2}}v_{0}^{-} - \partial_{x_{2}}v_{0}^{+} + \delta_{\tau,0}(1-h)\theta(x_{1})P(x_{1}) v_{0}^{+})(x_{1}, 0) = (1-h)\theta(x_{1})Q(x_{1}), (x_{1}, 0) \in I_{0}, \\ \partial_{x_{2}}v_{0}^{-}(x_{1}, -l) = 0, & (x_{1}, -l) \in I_{l}, \end{cases}$$
(8)

which is called homogenized problem for the problem (1). Here

$$I_l = \{x: x_2 = -l, x_1 \in (0, a)\};$$

 $\delta_{lpha,k}$ is the Kroneker symbol;

$$P(x_1) = \mathbb{E}\left(p(\xi_1, \omega) \sqrt{1 + (\theta(x_1) \partial_{\xi_1} F(\xi_1, \omega))^2}\right) = \mathbb{E}\left(\mathbf{p}(\omega) \sqrt{1 + (\theta(x_1) \partial_{\omega} \mathbf{F}(\omega))^2}\right),$$
$$Q(x_1) = \mathbb{E}\left(q(\xi_1, \omega) \sqrt{1 + (\theta(x_1) \partial_{\xi_1} F(\xi_1, \omega))^2}\right) = \mathbb{E}\left(\mathbf{q}(\omega) \sqrt{1 + (\theta(x_1) \partial_{\omega} \mathbf{F}(\omega))^2}\right).$$

Moreover the convergence of energy

$$E_{\varepsilon}(u_{\varepsilon}) := \int_{D_{\varepsilon}} |\nabla_{x}u_{\varepsilon}|^{2} dx + \varepsilon^{\tau} \int_{\Gamma_{\varepsilon}} \theta(x_{1}) p\left(\frac{x_{1}}{\varepsilon},\omega\right) u_{\varepsilon}^{2} d\sigma_{x} + \varepsilon^{\mu} k \int_{\Upsilon_{\varepsilon}} u_{\varepsilon}^{2} d\sigma_{x} \longrightarrow \int_{D_{0}} |\nabla v_{0}^{+}|^{2} dx + \int_{D_{2}} (h|\partial_{x_{2}}v_{0}^{-}|^{2} + 2\delta_{\mu,1}k|v_{0}^{-}|^{2}) dx + \delta_{\tau,0}(1-h) \int_{I_{0}} \theta(x_{1})P(x_{1})|v_{0}^{+}(x_{1},0)|^{2} dx_{1} =: E_{0}(v_{0})$$

$$(9)$$

holds true as $\varepsilon \to 0$ for almost all ω .

Theorem 2 (The case au < 0 and $\mu \geq 1$). For solutions $u_{arepsilon}$ to the problem (1) the limits

as $\varepsilon \to 0$ are valid for almost all ω , where the functions v_0^+ and v_0^- are respectively the solutions to the following problems:

$$\begin{cases}
-\Delta_{x} v_{0}^{+}(x) = f_{0}(x), & x \in D_{0} \\
v_{0}^{+}(x) = 0, & x \in \Gamma_{1} \cup I_{0} \\
\partial_{\nu} v_{0}^{+}(x) = 0, & x \in \partial D_{0} \setminus (\Gamma_{1} \cup I_{0}),
\end{cases}$$
(11)

$$\begin{cases} -h\partial_{x_{2}x_{2}}^{2}v_{0}^{-}(x) + 2\,\delta_{\mu,1}\,k\,v_{0}^{-}(x) &= h\,f_{0}(x) + \delta_{\beta,1}\,g_{0}(x), \quad x \in D_{2}, \\ v_{0}^{-}(x_{1},0) &= 0, \qquad (x_{1},0) \in I_{0}, \qquad (12) \\ \partial_{x_{2}}v_{0}^{-}(x_{1},-l) &= 0, \qquad (x_{1},-l) \in I_{l}, \end{cases}$$

which together are called the homogenized problem for the problem (1).

Moreover the convergence of the energy integrals

$$E_{\varepsilon}(u_{\varepsilon}) \rightarrow \int_{D_0} |\nabla v_0^+|^2 dx + h \int_{D_2} |\partial_{x_2} v_0^-|^2 dx + 2\delta_{\mu,1} k \int_{D_2} |v_0^-|^2 dx =: E_0(v_0^+) + E_0(v_0^-).$$
(13)

holds true as $\varepsilon \to 0$ for almost all ω .

Theorem 3 (The case $\mu < 1$). For the solution u_{ε} to the problem (1) for almost all ω the limits

$$\begin{array}{cccc} u_{\varepsilon} & \rightharpoonup & v_0^+ & \text{ in } H^1(D_0, \Gamma_1), \\ \widetilde{u_{\varepsilon}} & \rightarrow & 0 & \text{ in } L^2(D_2), \end{array} \right\} \quad \text{as } \varepsilon \to 0$$
 (14)

hold true, where the function v_0^+ is the solution to the problem

$$\begin{aligned}
-\Delta_x v_0^+(x) &= f_0(x), \quad x \in D_0 \\
v_0^+(x) &= 0, \quad x \in \Gamma_1 \cup I_0 \\
\partial_\nu v_0^+(x) &= 0, \quad x \in \partial D_0 \setminus (\Gamma_1 \cup I_0).
\end{aligned}$$
(15)

Moreover, for almost all ω the following convergence

$$E_{\varepsilon}(u_{\varepsilon}) \to \int_{D_0} |\nabla v_0^+|^2 \, dx =: E_0(v_0^+) \tag{16}$$

is valid as $\varepsilon \to 0$.

Auxiliary Lemmas

Lemma 1. Let $H(\xi_1, \omega)$ be a random statistically homogeneous function, such that $\|\mathbf{H}\|_{L^{\infty}(\Omega, \mu)} < \infty$ and

$$\mathbb{E}(H(\xi_1,\omega)) \equiv 0. \tag{17}$$

Then a.s.

$$\int_{I_0} H(\frac{x_1}{\varepsilon}, \omega) u(x_1) v(x_1) \, dx_1 \longrightarrow 0 \tag{18}$$

as $\varepsilon \to 0$ for any functions $u, v \in H^{\frac{1}{2}}(I_0)$.

Lemma 2. For any $u, v \in H^1(D_{\varepsilon})$ the following limit relations

$$\left| \int_{\Gamma_{\varepsilon}} \theta(x_1) q\left(\frac{x_1}{\varepsilon}, \omega\right) v(x) \ d\sigma_x - (1-h) \int_{I_0} \theta(x_1) Q(x_1) v(x_1, 0) \ dx_1 \right| \to 0, \tag{19}$$

$$\left| \int_{\Gamma_{\varepsilon}} \theta(x_1) p\left(\frac{x_1}{\varepsilon}, \omega\right) v(x) u(x) \, d\sigma_x - (1-h) \int_{I_0} \theta(x_1) P(x_1) v(x_1, 0) u(x_1, 0) \, dx_1 \right| \to 0$$
(20)

hold as $\varepsilon \to 0$ for almost all ω .

Auxiliary Lemmas

Boundary value problems in dense junctions with different nonhomogeneous conditions on the boundary of thin subdomains have specific difficulties. To homogenize problems in such junctions we use special integral identities. In this case the identity has the following form:

$$\frac{\varepsilon h}{2} \int_{\widehat{S}_{\varepsilon}} v \, dx_2 = \int_{\widehat{G}_{\varepsilon}} v \, dx - \varepsilon \int_{\widehat{G}_{\varepsilon}} Y_2\left(\frac{x_1}{\varepsilon}\right) \partial_{x_1} v \, dx, \qquad \forall v \in H^1(\widehat{G}_{\varepsilon}).$$
(21)

Here $Y_2(\xi) = -\xi + [\xi] + \frac{1}{2}$, where $[\xi]$ is the integer part of ξ ; S_{ε} is the union of the lateral sides of the rectangles G_{ε} .

Keeping in mind that $\max_{\mathbb{R}} |Y_2| \leq 1$, we get the inequality

$$\|v\|_{L^{2}(S_{\varepsilon})} \leq C_{2} \varepsilon^{-\frac{1}{2}} \|v\|_{H^{1}(G_{\varepsilon})}.$$
(22)

Using the standard approach we obtain

$$\|v\|_{L^{2}(B_{\varepsilon})} \leq C_{3} \|v\|_{H^{1}(G_{\varepsilon})},$$
(23)

where $B_{\varepsilon} = \Upsilon_{\varepsilon} \setminus S_{\varepsilon}$.

Comments

For 3D model with variable cross section of the rods we change the function Y_2 . Consider the following identity:

$$\varepsilon \int_{S_{\varepsilon}} \frac{\varphi(x) \, d\sigma_x}{\sqrt{1 + \varepsilon^2 |\varrho'(x_3)|^2}} = \int_{G_{\varepsilon}} \frac{l_{\omega}(x_3)}{|\omega(x_3)|} \, \varphi \, dx + \varepsilon \int_{G_{\varepsilon}} \nabla_{\xi'} Y_2(\xi', x_3) \big|_{\xi' = \frac{x'}{\varepsilon}} \cdot \nabla_{x'} \varphi \, dx \tag{24}$$

for any $\varphi \in H^1(G_{\varepsilon})$. Here Y_2 is 1-periodic in ξ_1 and ξ_2 function which satisfies

$$\begin{cases} \Delta_{\xi'} Y_2(\xi', x_3) &= \frac{l_{\omega}(x_3)}{|\omega(x_3)|} & \text{in } \omega(x_3), \\ \partial_{\nu'(\xi')} Y_2 &= 1 & \text{on } \partial\omega(x_3), \\ \int_{\omega(x_3)} Y_2(\xi', x_3) \, d\xi' &= 0, \end{cases}$$

$$(25)$$

where $\xi' = (\xi_1, \xi_2), \ \nu'(\xi') = (\nu_1(\xi'), \nu_2(\xi'))$ is outer normal to *D*.