
Orbifold deconstruction

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Introduction

Common formulations of dynamics (EOM, variational principles, ...) usually require (quasi-)trivial topology $\rightsquigarrow$ realize configuration space as either a submanifold or a quotient of some nice geometric structure: constrained dynamics vs **gauge symmetries**.

Dynamics formulated on covering space $\rightsquigarrow$ state identifications.

covering (deck) transformations = gauge transformations

New superselection sectors from 'twisted' boundary conditions.

Difficult computations (mostly numerical, e.g. lattice QCD).

Best understood for 2D conformal models (orbifolding).

Basic problems: structure of twisted sectors and fixed point resolution
(usually require *ad hoc* techniques).

Under control only in special cases:

- toroidal orbifolds (compactified free bosons, i.e. lattice models);
- holomorphic orbifolds (self-dual models);
- permutation orbifolds (permutation symmetries).

Moore's 'conjecture': all rational conformal models are GKO cosets or orbifolds thereof.

Can one reverse orbifolding?

Orbifolds

2D conformal model characterized by **chiral symmetry algebra** $\mathbb{V}$ (nice VOA) and **L-R coupling** (partition functions).

Primary fields (simple $\mathbb{V}$ -modules, aka. superselection sectors) characterized by their **conformal weight** h_p (lowest eigenvalue of L_0) and **chiral character** (trace function)

$$\chi_p(q) = \text{Tr}_p \left(q^{L_0 - c/24} \right)$$

describing the spectrum of L_0 .

Fusion rules: composition of superselection sectors (tensor products).

Fusion rules related to modular properties *via* **Verlinde's formula**.

For $G < \text{Aut}(\mathbb{V})$, the **chiral algebra** of the G -orbifold is the fixed point subalgebra $\mathbb{V}^G = \{v \in \mathbb{V} \mid gv = v \text{ for all } g \in G\}$.

Primaries of the G -orbifold from G -**twisted modules** of $\mathbb{V}$.

G permutes the G -twisted modules (**outer action**), with $h \in G$ taking a g -twisted module to a hgh^{-1} -twisted module $\rightsquigarrow$ **G -orbits organized into twisted sectors** labeled by conjugacy classes.

Twisted modules in the same G -orbit identified with each other (**'state identification'**).

Stabilizer $G_M = \{g \in G \mid gM = M\}$ of the twisted module M **represented projectively** on M (with associated 2-cocycle ϑ_M) $\rightsquigarrow$ M splits into isotypic components M_ϕ labeled by irreps $\phi \in \text{Irr}(G_M | \vartheta_M)$.

isotypic components $\longleftrightarrow$ primaries of the orbifold

G -orbits of (twisted) modules $\longleftrightarrow$ **blocks** of primaries

$$\text{number of primaries} = \frac{1}{|G|} \sum_{xy=yx} \sum_{M \in \text{Fix}(x,y)} \frac{\vartheta_M(x,y)}{\vartheta_M(y,x)}$$

$$\text{number of blocks} = \sum_M \frac{1}{[G : G_M]}$$

Each block $\mathfrak{b}$ characterized by **inertia subgroup** $\mathbf{I}_{\mathfrak{b}}$ (stabilizer G_M of any module M in the orbit corresponding to $\mathfrak{b}$) and 2-cocycle $\vartheta_{\mathfrak{b}} \in Z^2(\mathbf{I}_{\mathfrak{b}}, \mathbb{C})$.

Integrally spaced L_0 spectrum for untwisted modules $\rightsquigarrow$ **the conformal weights of primaries from a block in the untwisted sector differ by integers.**

Vacuum block $\mathfrak{b}_0$ (untwisted sector) has trivial cocycle $\rightsquigarrow$ all elements of $\mathfrak{b}_0$ have integer conformal weights, and correspond to (ordinary) irreps of G , with matching fusion rules and (quantum) dimensions.

The vacuum block is a **twister**: a set of primaries with integer conformal weights (and quantum dimensions) closed under fusion.

Question: can we identify the orbifold from its vacuum block?

Braiding restricted to a twister is involutive $\rightsquigarrow$ elements of the twister $\mathfrak{g}$ are the simple objects of a symmetric monoidal category.

Deligne's theorem: the subring of the fusion ring generated by a twister is isomorphic to the character ring of some finite group.

Fusion matrices

For a primary p define the fusion matrix

$$[\mathbf{N}(p)]_{qr} = N_{pq}^r$$

Span a commutative matrix algebra (the **Verlinde algebra** $\mathcal{V}$)

$$\mathbf{N}(p) \mathbf{N}(q) = \sum_r N_{pq}^r \mathbf{N}(r)$$

Commuting matrices with non-negative elements $\rightsquigarrow$ **common Perron-**

Frobenius eigenvector (**quantum dimension**) $d_p \geq 1$.

$$\sum_r N_{pq}^r d_r = d_p d_q$$

Modular S -matrix

$$S_{pq} = \frac{1}{\sqrt{\sum_r d_r^2}} \sum_r N_{pq}^r d_r \exp\{2\pi i(\mathbf{h}_p + \mathbf{h}_q - \mathbf{h}_r)\}$$

Verlinde's theorem:

$$\rho_p(\mathbf{N}(q)) = \frac{S_{qp}}{S_{0p}}$$

is an irrep of $\mathcal{V}$ for each primary p , where 0 denotes the **vacuum** primary.

irreducible representations of $\mathcal{V} \longleftrightarrow$ primary fields

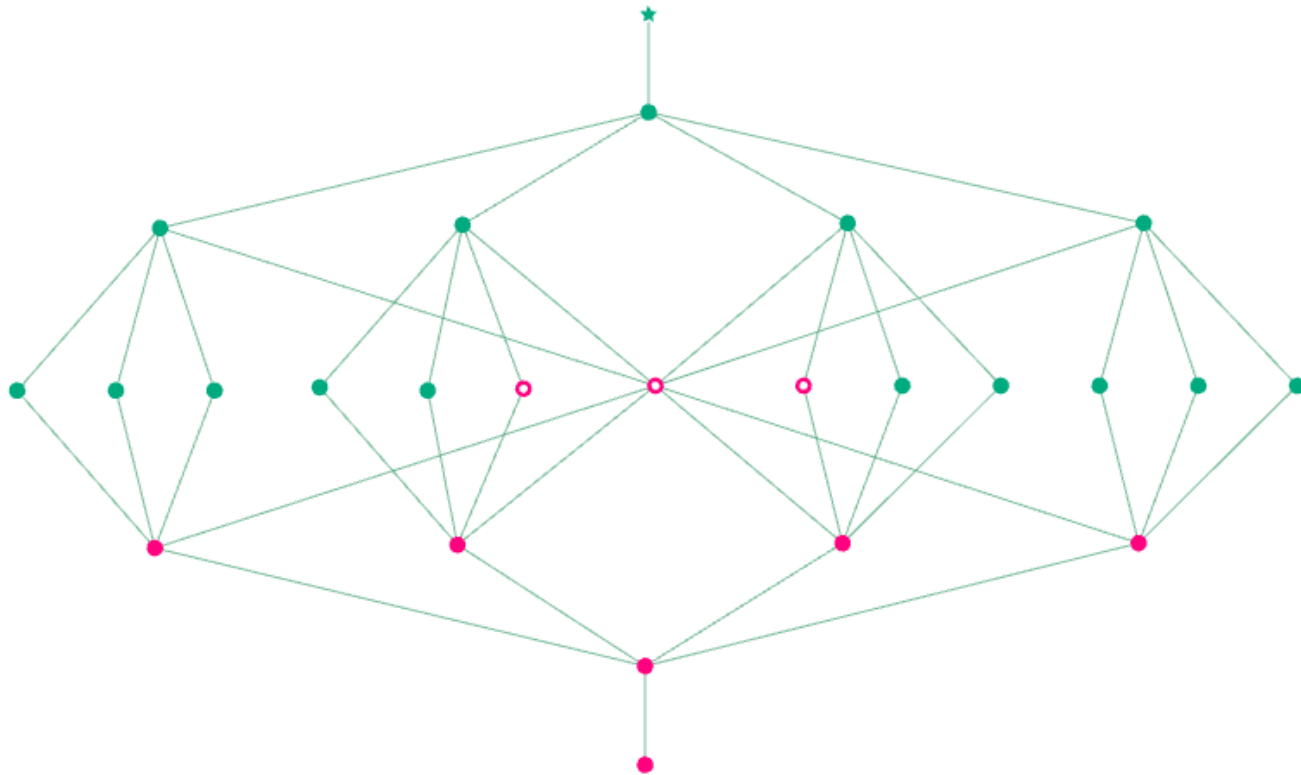
Perron's theorem $\rightsquigarrow$

$$|\rho_p(\mathbf{N}(q))| \leq \rho_0(\mathbf{N}(q)) = d_q$$

Twisters

A set $\mathfrak{g}$ of primaries is **fusion closed** if $N_{pq}^r > 0$ for $p, q \in \mathfrak{g}$ implies $r \in \mathfrak{g}$.

Form a **modular** (even Arguesian) **lattice** $\mathcal{L}$.



$\widehat{\mathfrak{g}}$: subalgebra of $\mathcal{V}$ spanned by the fusion matrices $N(\alpha)$ for $\alpha \in \mathfrak{g}$.

Primaries p and q belong to the same **twist class** if the restrictions to $\widehat{\mathfrak{g}}$ of the irreps ρ_p and ρ_q coincide.

twist classes $\leftrightarrow$ irreducible representations of $\widehat{\mathfrak{g}}$

$\rightsquigarrow$ the number of twist classes equals the cardinality of $\mathfrak{g}$.

Notation: for a class $\mathbb{C}$ and $\alpha \in \mathfrak{g}$ let $\alpha(\mathbb{C}) = \rho_p(\alpha)$ for any $p \in \mathbb{C}$.

Orthogonality relations: for $\alpha, \beta \in \mathfrak{g}$

$$\sum_{\mathbb{C}} \frac{\alpha(\mathbb{C}) \overline{\beta(\mathbb{C})}}{\|\mathbb{C}\|} = \begin{cases} 1 & \text{if } \alpha = \beta; \\ 0 & \text{otherwise} \end{cases}$$

where

$$\|\mathbb{C}\| = \frac{1}{\sum_{p \in \mathbb{C}} S_{0p}^2}$$

Second orthogonality: for twist classes $\mathbf{C}_1$ and $\mathbf{C}_2$

$$\sum_{\alpha \in \mathfrak{g}} \alpha(\mathbf{C}_1) \overline{\alpha(\mathbf{C}_2)} = \begin{cases} \|\mathbf{C}_1\| & \text{if } \mathbf{C}_1 = \mathbf{C}_2; \\ 0 & \text{otherwise.} \end{cases}$$

Trivial class $\mathfrak{g}^\perp$: twist class containing the vacuum primary $\mathbf{o}$.

Remark: $p \in \mathfrak{g}^\perp$ iff $\mathbf{h}_q - \mathbf{h}_p - \mathbf{h}_\alpha \in \mathbb{Z}$ whenever $N_{\alpha p}^q > 0$ for some $\alpha \in \mathfrak{g}$.

Product rule: If $p \in \mathfrak{g}^\perp$ and $N_{pq}^r > 0$, then q and r belong to the same twist class.

Consequence: $\mathfrak{g}^\perp \in \mathcal{L}$ and $(\mathfrak{g}^\perp)^\perp = \mathfrak{g}$, hence the lattice $\mathcal{L}$ is self-dual.

Blocks of $\mathfrak{g}$ = twist classes of $\mathfrak{g}^\perp$ (number of blocks = cardinality of $\mathfrak{g}^\perp$).

Blocks partition the set of primaries: p and q belong to the same block iff $N_{\alpha p}^q > 0$ for some $\alpha \in \mathfrak{g}$ ($\mathfrak{g}$ is the trivial block).

Restriction of fusion matrices to the elements of a block $\mathfrak{b}$ provides an integral representation $N_{\mathfrak{b}}$ of $\widehat{\mathfrak{g}}$.

$\left. \begin{array}{l} \text{classes: irreducible} \\ \text{blocks: integral} \end{array} \right\} \text{representations of } \widehat{\mathfrak{g}}$

Overlap $\langle \mathfrak{b}, \mathbf{C} \rangle$: multiplicity of irrep $\rho_{\mathbf{C}}$ in $N_{\mathfrak{b}}$ (non-negative integer)

$$\langle \mathfrak{b}, \mathbf{C} \rangle = \sum_{p \in \mathfrak{b}} \sum_{q \in \mathbf{C}} |S_{pq}|^2$$

$$|\mathfrak{b}| = \sum_{\mathbf{C}} \langle \mathfrak{b}, \mathbf{C} \rangle \text{ and } |\mathbf{C}| = \sum_{\mathfrak{b}} \langle \mathfrak{b}, \mathbf{C} \rangle.$$

Note: $\langle \mathfrak{b}, \mathbf{C} \rangle = 0$ implies $S_{pq} = 0$ for all $p \in \mathfrak{b}$ and $q \in \mathbf{C}$.

$$\langle \mathfrak{g}, \mathbf{C} \rangle = 1 \text{ for every class } \mathbf{C} \text{ and } \langle \mathfrak{b}, \mathfrak{g}^{\perp} \rangle = 1 \text{ for every block } \mathfrak{b}.$$

Integrality theorem: $\mathfrak{g} \subseteq \mathfrak{g}^\perp$ implies $\mathbf{d}_\alpha \in \mathbb{Z}$ and $\mathbf{h}_\alpha \in \frac{1}{2}\mathbb{Z}$ for all $\alpha \in \mathfrak{g}$.

$\mathfrak{g} \subseteq \mathfrak{g}^\perp$ **iff** $N_{\alpha\beta}^\gamma > 0$ for $\alpha, \beta, \gamma \in \mathfrak{g}$ implies $\mathbf{h}_\gamma - \mathbf{h}_\alpha - \mathbf{h}_\beta \in \mathbb{Z}$ **iff** every class is a union of blocks.

$\mathfrak{g} \in \mathcal{L}$ is a **twister** if $\mathbf{h}_\alpha \in \mathbb{Z}$ for all $\alpha \in \mathfrak{g}$ (consequently $\mathfrak{g} \subseteq \mathfrak{g}^\perp$).

Every $\mathfrak{g} \subseteq \mathfrak{g}^\perp$ is a twister or a $\mathbb{Z}_2$ -extension thereof ('**swister**'?).

If $\mathfrak{g} \in \mathcal{L}$ is a twister, then

1. a block is contained in the trivial class $\mathfrak{g}^\perp$ iff the conformal weights of its elements differ by integers;
2. to each block $\mathfrak{b}$ corresponds an algebraic integer $\mathcal{D}_\mathfrak{b}$ ('**block dimension**') such that $\frac{\mathbf{d}_p}{\mathcal{D}_\mathfrak{b}} \in \mathbb{Z}_+$ for all $p \in \mathfrak{b}$.

A twister determines a Tannakian subcategory of the modular tensor category associated to the conformal model $\rightsquigarrow$

the ring $\widehat{\mathfrak{g}}$ associated to a twister is isomorphic to the representation ring of some finite group G (Deligne's theorem).

Problem: there may exist several nonisomorphic groups with isomorphic representation rings (e.g. $\mathbb{D}_8$, the dihedral group of order 8, and Q , the group of unit quaternions). How to distinguish between the different possibilities?

Solution: use the λ -ring structure (exterior powers) of the representation ring, i.e. the Adams operations.

Adams operations

For $n \in \mathbb{Z}$ and elements $\alpha, \beta \in \mathfrak{g}$ of a twister $\mathfrak{g}$ let

$$Z_n(\alpha, \beta) = \sum_{p,q} N_{\alpha p}^q S_{\beta p} S_{0q} e^{2\pi i n(\mathbf{h}_\alpha + \mathbf{h}_p - \mathbf{h}_q)}$$

There exists an endomorphism $\Psi^n \in \text{End}(\widehat{\mathfrak{g}})$ such that

- $\Psi^n(\alpha) = \sum_{\beta \in \mathfrak{g}} Z_n(\alpha, \beta) \beta$ for $\alpha \in \mathfrak{g}$;
- $\Psi^n(\alpha) = \alpha^n$ for invertible elements $\alpha \in \widehat{\mathfrak{g}}$;
- Ψ^n is functorial;
- $\Psi^n \circ \Psi^m = \Psi^{nm}$.

$\Psi^n \in \text{End}(\widehat{\mathfrak{g}})$ is the n^{th} **Adams operation** on $\widehat{\mathfrak{g}}$, endowing the latter with the structure of a **λ -ring**, able to distinguish (but for **Brauer pairs**) non-isomorphic groups with identical fusion rules, e.g. $\mathbb{D}_8$ and Q .

Ψ^n permutes the characteristic functions of the twist classes $\rightsquigarrow$ existence of power maps $\mathcal{C} \mapsto \mathcal{C}^n$ for each $n \in \mathbb{Z}$.

Order of the twist class $\mathcal{C}$ = least positive integer n such that $\mathcal{C}^n = \mathfrak{g}^\perp$.

If a block $\mathfrak{b}$ is contained in a twist class of order n , then the conformal weights inside $\mathfrak{b}$ can only differ by integer multiples of $\frac{1}{n}$.

Orbifold deconstruction

vacuum block	twister
irrep of twist group	element of twister
twisted sectors	twist classes
orbits of twisted modules	blocks
stabilizer of orbit	inertia subgroup

Step 1: select a twister.

Step 2: determine the corresponding twist classes and blocks.

Step 3: compute the character table and the power maps.

Step 4: identify the twist group G (up to possible Brauer pairs).

Step 5: for each block $\mathfrak{b}$ contained in the trivial class determine the inertia subgroup $\mathbf{I}_{\mathfrak{b}}$ and associated 2-cocycle $\vartheta_{\mathfrak{b}} \in Z^2(\mathbf{I}_{\mathfrak{b}}, \mathbb{C}^\times)$.

Remark. For most applications it is enough to know the order $\mathbf{e}_{\mathfrak{b}}$ of (the cohomology class of) $\vartheta_{\mathfrak{b}}$ and the index $[G:\mathbf{I}_{\mathfrak{b}}]$ of the inertia subgroup.

Step 6: compute the spectrum of the deconstructed model.

Remark. To each block $\mathfrak{b} \subseteq \mathfrak{g}^\perp$ correspond $[G:\mathbf{I}_{\mathfrak{b}}]$ different primaries of respective conformal weights $\mathbf{h}_{\mathfrak{b}} = \min \{\mathbf{h}_p \mid p \in \mathfrak{b}\}$, quantum dimensions

$$\mathbf{d}_{\mathfrak{b}} = \frac{\mathcal{D}_{\mathfrak{b}}}{\mathbf{e}_{\mathfrak{b}} [G:\mathbf{I}_{\mathfrak{b}}]}$$

and chiral characters

$$\chi_{\mathfrak{b}}(\tau) = \frac{e_{\mathfrak{b}}}{\mathcal{D}_{\mathfrak{b}}} \sum_{p \in \mathfrak{b}} d_p \chi_p(\tau)$$

Step 7: if the above did not single out the deconstructed model, compute the block-fusion coefficients

$$\mathcal{N}_{\mathfrak{a}\mathfrak{b}}^{\mathfrak{c}} = \frac{e_{\mathfrak{a}} e_{\mathfrak{b}} e_{\mathfrak{c}}}{|\mathcal{I}_{\mathfrak{a}}| |\mathcal{I}_{\mathfrak{b}}|} \sum_{p \in \mathfrak{a}} \sum_{q \in \mathfrak{b}} \sum_{r \in \mathfrak{c}} N_{pq}^r \frac{d_p d_q d_r}{\mathcal{D}_{\mathfrak{a}} \mathcal{D}_{\mathfrak{b}} \mathcal{D}_{\mathfrak{c}}} \quad (1)$$

that characterize the fusion rules of the deconstructed model (but for fixed-point resolution).

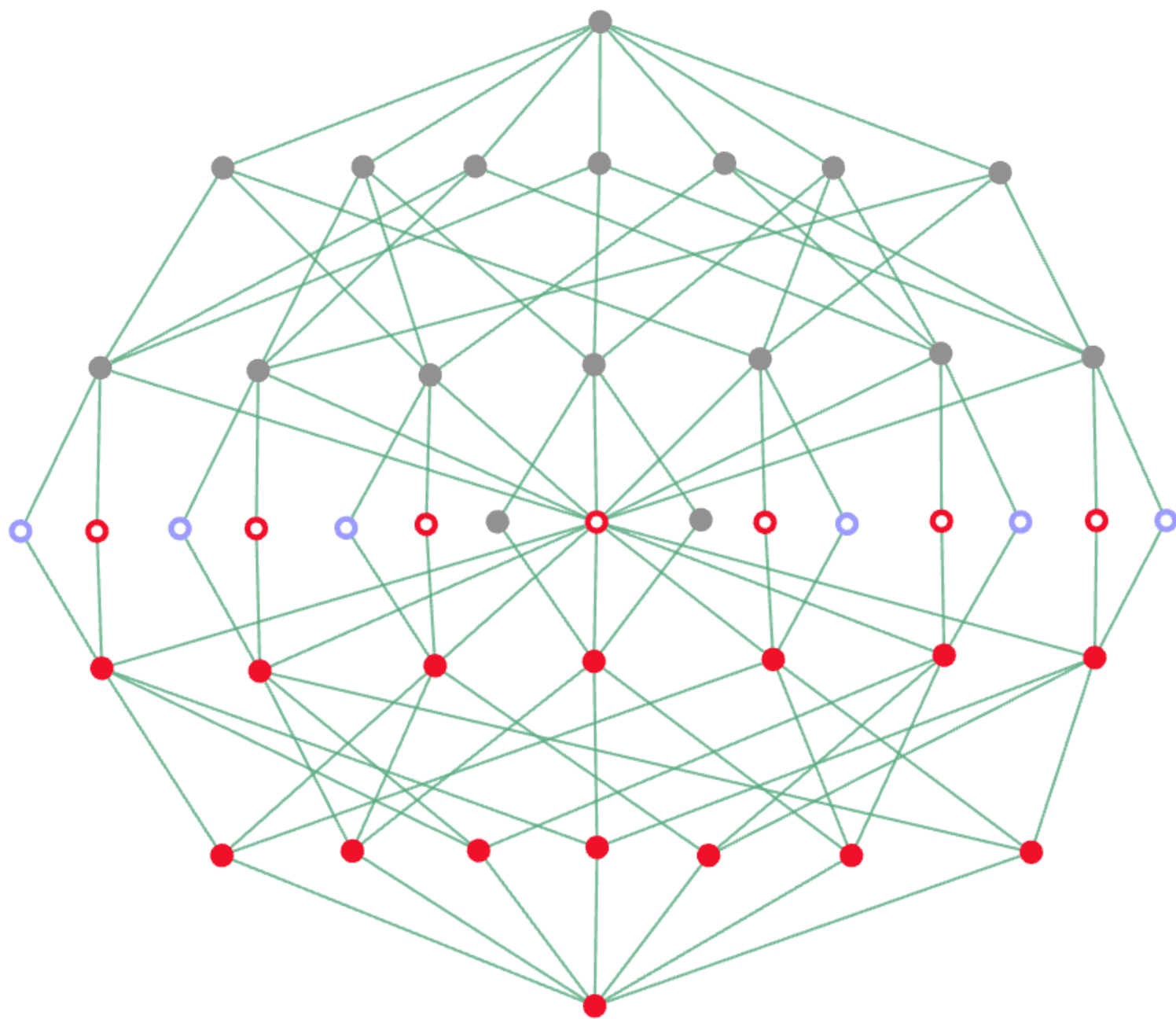
A permutation orbifold example

$\mathbf{D}$ = transitive permutation group of degree 4 generated by the cyclic permutations $(1, 2, 3, 4)$ and $(2, 4)$ (isomorphic to the dihedral group of order 8, the symmetry group of a square).

Permutation orbifold $\mathbb{V}^{\mathbf{D}}$: conformal model of central charge $c = 96$, with 22 primaries of known conformal weights, chiral characters, etc.

Moonshine module is self-dual $\rightsquigarrow$ fusion rules and modular properties described by the (untwisted) double of $\mathbf{D}$.

45 fusion closed sets (of which 22 twistors), with Hasse-diagram



7 maximal twisters, all self-dual $\rightsquigarrow$ corresponding deconstructions are self-dual too (with only one primary).

twist group	trace function
$\mathbb{D}_8$	$J(q)^4$
$\mathbb{D}_8$	$J(q)^4 - 590652J(q)^2 - 64481280J(q) + 55552950252$
$\mathbb{D}_8$	$J(q)^4 - 590652J(q)^2 - 64481280J(q) + 55552359600$
$\mathbb{D}_8$	$J(q)^4 - 590652J(q)^2 - 64481277J(q) + 55552950252$
$\mathbb{D}_8$	$J(q)^4 - 393768J(q)^2 + 38763309456 = (J(q)^2 - 196884)^2$
$\mathbb{D}_8$	$J(q)^4 - 393768J(q)^2 - 42987519J(q) - 1728009288$
$\mathbb{Z}_2^3$	$J(q)^4 - 393768J(q)^2 - 42987519J(q) + 37035300168$

Table 1. Maximal deconstructions of $\mathbb{V}^{\mathfrak{d}} \wr \mathbf{D}$

Maximal deconstructions have different trace functions $\rightsquigarrow$ they all differ from each other.

Lessons:

1. one and the same conformal model can have several different maximal deconstructions, possibly with non-isomorphic twist groups;
2. one needs to know the power maps in order to identify the twist group;
3. maximal deconstructions of a permutation orbifold related to **replication identities**, e.g. (from deconstruction no. 7)

$$\begin{aligned} & J(2\tau)^2 + J\left(\frac{\tau}{2}\right)^2 + J\left(\frac{\tau+1}{2}\right)^2 \\ &= J(\tau)^4 - 787536J(\tau)^2 - 85975038J(\tau) + 74070600336 \end{aligned}$$

Addendum: $L\left(\frac{1}{2}, 0\right) \wr \mathbb{S}_3$ (aka. $\mathcal{F}(3)^{\mathbb{S}_3}$, cf. Penn's talk) has 49 primaries, but only 3 (non-trivial) twistors. The corresponding decompositions are

twist group	deconstructed model
$\mathbb{Z}_2$	$L\left(\frac{1}{2}, 0\right) \wr \mathbb{Z}_3$
$\mathbb{S}_3$	$L\left(\frac{1}{2}, 0\right)^{\otimes 3}$
$\mathbb{S}_4$	$\mathrm{SU}(2)_2$

Each deconstruction is $N=1$ superconformal, with $L\left(\frac{1}{2}, 0\right) \wr \mathbb{Z}_3$ and (the trivial) $L\left(\frac{1}{2}, 0\right) \wr \mathbb{S}_3$ isolated in the moduli space of $c=\frac{3}{2}$ superconformal models (Cappelli and d'Appollonio, JHEP0208:039, 2002).

Summary

- effective procedure for orbifold deconstruction
- lattice structure of fusion closed sets
- Adams-operations for twistors
- integrality of quantum dimensions
- structure of twisted modules

and open questions

- characterization of primitive models?
- Moore's conjecture?
- relation between different maximal deconstructions?
- generalized deconstruction using swisters?
- Brauer characters?