# THE CAUCHY PROBLEM FOR ONE-DIMENSIONAL FLOW OF A COMPRESSIBLE VISCOUS FLUID: STABILIZATION OF THE SOLUTION 

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#### Abstract

We analyze the Cauchy problem for non-stationary 1-D flow of a compressible viscous and heat-conducting fluid, assuming that it is in the thermodynamical sense perfect and polytropic. This problem has a unique generalized solution on $\mathbb{R} \times] 0, T$ for each $T>0$. Supposing that the initial functions are small perturbations of the constants and using some a priori estimates for the solution independent of $T$, we prove a stabilization of the solution.


## 1. Introduction

In this paper we analyze the Cauchy problem for non-stationary 1-D flow of a compressible viscous and heat-conducting fluid. It is assumed that the fluid is thermodynamically perfect and polytropic. The same model has been mentioned in [1], where a global-in-time existence theorem for generalized solution is given without a rigorous proof. Here, we approach to our problem as the special case of the Cauchy problem for a micropolar fluid that is considered in [5] and [6]. Therefore we know that this problem has a unique generalized solution on $\mathbb{R} \times] 0, T[$, for each $T>0$ and that the mass density and temperature are strictly positive.

Assuming that the initial functions are small perturbations of the constants, we first derive a priori estimates for a solution independent of $T$ and then analyze the behavior of the solution as $T \rightarrow \infty$. We use some ideas of Ya. I. Kanel' ([4]) applied to the case of stabilization of Hölder continuous solution for the same model.

Key words and phrases. Compressible viscous fluid, the Cauchy problem, stabilization.

## 2. Statement of the problem and the main result

Let $\rho, v$ and $\theta$ denote, respectively, the mass density, velocity, and temperature of the fluid in the Lagrangean description. Supposing that in [5] the microrotation is equal to zero, we obtain the problem considered as follows:

$$
\begin{gather*}
\frac{\partial \rho}{\partial t}+\rho^{2} \frac{\partial v}{\partial x}=0  \tag{2.1}\\
\frac{\partial v}{\partial t}=\frac{\partial}{\partial x}\left(\rho \frac{\partial v}{\partial x}\right)-K \frac{\partial}{\partial x}(\rho \theta)  \tag{2.2}\\
\frac{\partial \theta}{\partial t}=-K \rho \theta \frac{\partial v}{\partial x}+\rho\left(\frac{\partial v}{\partial x}\right)^{2}+D \frac{\partial}{\partial x}\left(\rho \frac{\partial \theta}{\partial x}\right) \tag{2.3}
\end{gather*}
$$

in $\mathbb{R} \times \mathbb{R}^{+}$, where $K$ and $D$ are positive constants. The equations (2.1)-(2.3) are, respectively, local forms of the conservation laws for the mass, momentum and energy ${ }^{1}$. We take the following non-homogeneous initial conditions:

$$
\begin{gather*}
\rho(x, 0)=\rho_{0}(x),  \tag{2.4}\\
v(x, 0)=v_{0}(x),  \tag{2.5}\\
\theta(x, 0)=\theta_{0}(x) \tag{2.6}
\end{gather*}
$$

for $x \in \mathbb{R}$, where $\rho_{0}, v_{0}$ and $\theta_{0}$ are given functions. We assume that there exist the constants $m, M \in \mathbb{R}^{+}$, such that

$$
\begin{equation*}
m \leq \rho_{0}(x) \leq M, \quad m \leq \theta_{0}(x) \leq M, \quad x \in \mathbb{R} \tag{2.7}
\end{equation*}
$$

In the papers [5] and [6] it was proved that for

$$
\begin{equation*}
\rho_{0}-1, v_{0}, \theta_{0}-1 \in \mathrm{H}^{1}(\mathbb{R}) \tag{2.8}
\end{equation*}
$$

the problem (2.1)-(2.6) has, for each $T \in \mathbb{R}^{+}$, a unique generalized solution

$$
\begin{equation*}
(x, t) \mapsto(\rho, v, \theta)(x, t), \quad(x, t) \in \Pi=\mathbb{R} \times] 0, T[ \tag{2.9}
\end{equation*}
$$

with the properties:

$$
\begin{equation*}
\rho-1 \in \mathrm{~L}^{\infty}\left(0, T ; \mathrm{H}^{1}(\mathbb{R})\right) \cap \mathrm{H}^{1}(\Pi), \tag{2.10}
\end{equation*}
$$

$$
\begin{equation*}
v, \theta-1 \in \mathrm{~L}^{\infty}\left(0, T ; \mathrm{H}^{1}(\mathbb{R})\right) \cap \mathrm{H}^{1}(\Pi) \cap \mathrm{L}^{2}\left(0, T ; \mathrm{H}^{2}(\mathbb{R})\right) \tag{2.11}
\end{equation*}
$$

Using the results from [5] and [1] we can easily conclude that

$$
\begin{equation*}
\theta, \rho>0 \text { in } \Pi . \tag{2.12}
\end{equation*}
$$

[^0]We denote by $\mathrm{B}^{k}(\mathbb{R}), k \in \mathbb{N}_{0}$, the Banach space

$$
\begin{equation*}
\mathrm{B}^{k}(\mathbb{R})=\left\{u \in \mathrm{C}^{k}(\mathbb{R}): \lim _{|x| \rightarrow \infty}\left|D^{n} u(x)\right|=0,0 \leq n \leq k\right\} \tag{2.13}
\end{equation*}
$$

where $D^{n}$ is $n$-th derivative. The norm of the space $\mathrm{B}^{k}(\mathbb{R})$ is defined by

$$
\begin{equation*}
\|u\|_{\mathrm{B}^{k}(\mathbb{R})}=\sup _{n \leq k}\left\{\sup _{x \in \mathbb{R}}\left|D^{n} u(x)\right|\right\} \tag{2.14}
\end{equation*}
$$

From Sobolev's embedding theorem ([2, Chapter IV]) and the theory of vector-valued distributions ([3, pp. 467-480]) one can conclude that from (2.10) and (2.11) follows

$$
\begin{gather*}
\rho-1 \in \mathrm{~L}^{\infty}\left(0, T ; \mathrm{B}^{0}(\mathbb{R})\right) \cap \mathrm{C}\left([0, T] ; \mathrm{L}^{2}(\mathbb{R})\right),  \tag{2.15}\\
v, \theta-1 \in \mathrm{~L}^{2}\left(0, T ; \mathrm{B}^{1}(\mathbb{R})\right) \cap \mathrm{C}\left([0, T] ; \mathrm{H}^{1}(\mathbb{R})\right) \cap \mathrm{L}^{\infty}\left(0, T ; \mathrm{B}^{0}(\mathbb{R})\right) \tag{2.16}
\end{gather*}
$$

and hence

$$
\begin{equation*}
v, \theta-1 \in \mathrm{C}\left([0, T] ; \mathrm{B}^{0}(\mathbb{R})\right), \quad \rho \in \mathrm{L}^{\infty}(\Pi) \tag{2.17}
\end{equation*}
$$

From (2.7) and (2.8) it is easy to see that there exist the constants $E_{1}$, $E_{2}, E_{3}, M_{1} \in \mathbb{R}^{+}, M_{1}>1$, such that
(2.18) $\frac{1}{2} \int_{\mathbb{R}} v_{0}^{2} d x+K \int_{\mathbb{R}}\left(\frac{1}{\rho_{0}}-\ln \frac{1}{\rho_{0}}-1\right) d x+\int_{\mathbb{R}}\left(\theta_{0}-\ln \theta_{0}-1\right) d x=E_{1}$,

$$
\begin{equation*}
\frac{1}{2} \int_{\mathbb{R}} \frac{\left(\rho_{0}^{\prime}\right)^{2}}{\rho_{0}^{2}} d x+\int_{\mathbb{R}} v_{0}^{\prime} \ln \frac{1}{\rho_{0}} d x \leq E_{3} \tag{2.20}
\end{equation*}
$$

$$
\begin{equation*}
\frac{1}{2} \int_{\mathbb{R}}\left(\left(v_{0}^{\prime}\right)^{2}+\left(\theta_{0}^{\prime}\right)^{2}\right) d x=E_{2} \tag{2.19}
\end{equation*}
$$

$$
\begin{equation*}
\sup _{|x|<\infty} \theta_{0}(x)<M_{1} \tag{2.21}
\end{equation*}
$$

Suppose that the quantities $\underline{\eta}$ and $\bar{\eta}$ are such that $\underline{\eta}<0<\bar{\eta}$ and

$$
\begin{equation*}
\int_{\underline{\eta}}^{0} \sqrt{e^{\eta}-1-\eta} d \eta=\int_{0}^{\bar{\eta}} \sqrt{e^{\eta}-1-\eta} d \eta=E_{5} \tag{2.22}
\end{equation*}
$$

where

$$
\begin{equation*}
E_{5}=2 \sqrt{\frac{E_{1} E_{4}}{K}}, \quad E_{4}=2 \mu E_{1}\left(1+M_{1}+\frac{E_{3}}{E_{1}}\right), \quad \mu=\max \left\{\frac{K}{2 D}, 1\right\} \tag{2.23}
\end{equation*}
$$

Let

$$
\begin{equation*}
\underline{u}=\exp \underline{\eta}, \quad \bar{u}=\exp \bar{\eta} \tag{2.24}
\end{equation*}
$$

The aim of this work is to prove the following theorem.

THEOREM 2.1. Suppose that the initial functions satisfy (2.7), (2.8) and the following conditions:

$$
\begin{equation*}
E_{1} \int_{\mathbb{R}}\left(v_{0}^{\prime}\right)^{2} d x<\left(\frac{D}{16} \frac{u}{\bar{u}}\right)^{2} \tag{2.25}
\end{equation*}
$$

$$
\begin{align*}
& 2 E_{1}\left(E_{1} \frac{\bar{u}}{\underline{u}}\left(1+M_{1}\right) M_{1}\left(\left(16 E_{4}\right)^{2} \frac{\bar{u}}{\underline{u}}+\frac{K^{2}}{D} M_{1}\right)+2 K M_{1} E_{4}+E_{2}\right)  \tag{2.26}\\
& <\min \left\{\left(\frac{D}{16} \frac{u}{\bar{u}}\right)^{2},\left(\int_{1}^{M_{1}} \sqrt{s-1-\ln s} d s\right)^{2},\left(\int_{0}^{1} \sqrt{s-1-\ln s} d s\right)^{2}\right\}
\end{align*}
$$

then

$$
\begin{equation*}
\rho(x, t) \rightarrow 1, \quad v(x, t) \rightarrow 0, \quad \theta(x, t) \rightarrow 1, \text { when } t \rightarrow \infty \tag{2.27}
\end{equation*}
$$

uniformly with respect to all $x \in \mathbb{R}$.
Remark 2.2. Conditions (2.25) and (2.26) mean that $E_{1}, E_{2}$ and $E_{3}$ are sufficiently small. In other words the initial functions are small perturbations of the constants.

In the proof of Theorem 2.1. we apply some ideas of [4], where a stabilization of the solution that is Hölder continuous was proved for the same model.

## 3. A PRORI ESTIMATES FOR $\rho, v$ and $\theta$

Considering stabilization problem, one has to prove some a priori estimates for the solution independent of $T$, which is the main difficulty. Some of our considerations are similar to those of [4]. First we derive the energy equation for the solution of problem (2.1)-(2.3) under the conditions indicated above and we estimate the function $\rho^{-1}$.

Lemma 3.1. For each $t>0$ we have

$$
\begin{align*}
\frac{1}{2} \int_{\mathbb{R}} v^{2} d x+\int_{\mathbb{R}}(\theta & -\ln \theta-1) d x+K \int_{\mathbb{R}}\left(\frac{1}{\rho}-\ln \frac{1}{\rho}-1\right) d x  \tag{3.1}\\
& +\int_{0}^{t} \int_{\mathbb{R}}\left(\frac{\rho}{\theta}\left(\frac{\partial v}{\partial x}\right)^{2}+D \frac{\rho}{\theta^{2}}\left(\frac{\partial \theta}{\partial x}\right)^{2}\right) d x d \tau=E_{1}
\end{align*}
$$

where $E_{1}$ is defined by (2.18).
Proof. Multiplying (2.1), (2.2) and (2.3), respectively, by $K \rho^{-1}(1-$ $\left.\rho^{-1}\right), v$ and $1-\theta^{-1}$, integrating by parts over $\mathbb{R}$ and over $] 0, t[$ and taking into account (2.15) and (2.16), after addition of the obtained equations we easily get equality (3.1) independently of $t$.

Lemma 3.2. For each $t>0$ exist the strictly positive quantities $\underline{u}_{1}=$ $\underline{u}_{1}(\bar{\theta}(t))$ and $\bar{u}_{1}=\bar{u}_{1}(\bar{\theta}(t))$ such that

$$
\begin{equation*}
\left.\underline{u}_{1} \leq \rho^{-1}(x, \tau) \leq \bar{u}_{1}, \quad(x, \tau) \in \mathbb{R} \times\right] 0, t[ \tag{3.2}
\end{equation*}
$$

where

$$
\begin{equation*}
\bar{\theta}(t)=\sup _{(x, \tau) \in \mathbb{R} \times] 0, t[ } \theta(x, \tau) \tag{3.3}
\end{equation*}
$$

Proof. We multiply (2.2) by $\frac{\partial}{\partial x} \ln \left(\frac{1}{\rho}\right)$, integrate over $\mathbb{R}$ and $] 0, t[$ and use the following equalities, which follow using (2.1) and (2.15)-(2.17):

$$
\begin{align*}
& \int_{0}^{t} \int_{\mathbb{R}} \frac{\partial v}{\partial t} \frac{\partial}{\partial x}\left(\ln \frac{1}{\rho}\right) d x d \tau=-\int_{0}^{t} \int_{\mathbb{R}} \frac{\partial^{2} v}{\partial x \partial t} \ln \frac{1}{\rho} d x d \tau  \tag{3.5}\\
& =-\left.\int_{\mathbb{R}} \frac{\partial v}{\partial x} \ln \frac{1}{\rho} d x\right|_{0} ^{t}+\int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial v}{\partial x}\right)^{2} d x d \tau, \\
& -\left.\int_{\mathbb{R}} \frac{\partial v}{\partial x} \ln \frac{1}{\rho} d x\right|_{0} ^{t}=-\int_{\mathbb{R}} \frac{\partial}{\partial x} v \ln \frac{1}{\rho} d x+\int_{\mathbb{R}} v_{0}^{\prime} \ln \frac{1}{\rho_{0}} d x  \tag{3.6}\\
& =\int_{\mathbb{R}} v \frac{\partial}{\partial x}\left(\ln \frac{1}{\rho}\right) d x+\int_{\mathbb{R}} v_{0}^{\prime} \ln \frac{1}{\rho_{0}} d x, \tag{3.7}
\end{align*}
$$

$$
\begin{equation*}
\frac{\partial}{\partial x}\left(\rho \frac{\partial v}{\partial x}\right) \frac{\partial}{\partial x}\left(\ln \frac{1}{\rho}\right)=\frac{1}{2} \frac{\partial}{\partial t}\left(\frac{\partial}{\partial x}\left(\ln \frac{1}{\rho}\right)\right)^{2} \tag{3.4}
\end{equation*}
$$

We also take into account the following inequalities obtained by the Young's inequality

$$
\begin{align*}
& \rho^{2} \frac{\partial \theta}{\partial x} \frac{\partial}{\partial x}\left(\frac{1}{\rho}\right) \leq \frac{1}{2} \frac{\rho}{\theta}\left(\frac{\partial \theta}{\partial x}\right)^{2}+\frac{1}{2} \theta \rho^{3}\left(\frac{\partial}{\partial x}\left(\frac{1}{\rho}\right)\right)^{2}  \tag{3.8}\\
& v \frac{\partial}{\partial x} \ln \left(\frac{1}{\rho}\right)=v \rho \frac{\partial}{\partial x}\left(\frac{1}{\rho}\right) \leq \frac{1}{4} \rho^{2}\left(\frac{\partial}{\partial x}\left(\frac{1}{\rho}\right)\right)^{2}+v^{2} \tag{3.9}
\end{align*}
$$

After simple transformations we get

$$
\begin{align*}
& \frac{1}{4} \int_{\mathbb{R}} \rho^{2}\left(\frac{\partial}{\partial x}\left(\frac{1}{\rho}\right)\right)^{2} d x+\frac{K}{2} \int_{0}^{t} \int_{\mathbb{R}} \theta \rho^{3}\left(\frac{\partial}{\partial x}\left(\frac{1}{\rho}\right)\right)^{2} d x d \tau \\
& \leq \int_{\mathbb{R}} v^{2} d x+\frac{K}{2} \int_{0}^{t} \int_{\mathbb{R}} \frac{\rho}{\theta}\left(\frac{\partial \theta}{\partial x}\right)^{2} d x d \tau+\int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial v}{\partial x}\right)^{2} d x d \tau  \tag{3.10}\\
& \quad+\int_{\mathbb{R}} v_{0}^{\prime} \ln \left(\frac{1}{\rho_{0}}\right) d x+\frac{1}{2} \int_{\mathbb{R}} \frac{1}{\rho_{0}^{2}}\left(\rho_{0}^{\prime}\right)^{2} d x
\end{align*}
$$

for each $t>0$. Using (3.1) and (2.20) from (3.10) we obtain

$$
\begin{equation*}
\frac{1}{4} \int_{\mathbb{R}} \rho^{2}\left(\frac{\partial}{\partial x}\left(\frac{1}{\rho}\right)\right)^{2} d x+\frac{K}{2} \int_{0}^{t} \int_{\mathbb{R}} \theta \rho^{3}\left(\frac{\partial}{\partial x}\left(\frac{1}{\rho}\right)\right)^{2} d x d \tau \leq K_{1}(\bar{\theta}(t)) \tag{3.11}
\end{equation*}
$$

where

$$
\begin{equation*}
K_{1}(\bar{\theta}(t))=2 \mu E_{1}\left(1+\bar{\theta}(t)+\frac{E_{3}}{E_{1}}\right) \tag{3.12}
\end{equation*}
$$

and $\mu, E_{1}$ and $E_{3}$ are defined by (2.23), (2.18) and (2.20).
Now we define the increasing function $\psi$ by

$$
\begin{equation*}
\psi(\eta)=\int_{0}^{\eta} \sqrt{e^{\xi}-1-\xi} d \xi \tag{3.13}
\end{equation*}
$$

One can conclude the following

$$
\begin{equation*}
\left|\psi\left(\ln \frac{1}{\rho}\right)\right|=\left|\int_{0}^{\ln \frac{1}{\rho}} \psi^{\prime}(\xi) d \xi\right| \leq\left|\int_{\mathbb{R}} \psi^{\prime}\left(\ln \frac{1}{\rho}\right) \rho \frac{\partial}{\partial x}\left(\frac{1}{\rho}\right) d x\right| \tag{3.14}
\end{equation*}
$$

Using (3.1), (3.11) and the Hölder's inequality we get

$$
\begin{align*}
\left|\psi\left(\ln \frac{1}{\rho}\right)\right| & \leq\left(\int_{\mathbb{R}}\left(\frac{1}{\rho}-1-\ln \frac{1}{\rho}\right) d x\right)^{\frac{1}{2}}\left(\int_{\mathbb{R}} \rho^{2}\left(\frac{\partial}{\partial x}\left(\frac{1}{\rho}\right)\right)^{2} d x\right)^{\frac{1}{2}}  \tag{3.15}\\
& \leq K_{2}(\bar{\theta}(t))
\end{align*}
$$

where

$$
\begin{equation*}
K_{2}(\bar{\theta}(t))=2 E_{1}\left(\frac{2 \mu}{K}\left(1+\bar{\theta}(t)+\frac{E_{3}}{E_{1}}\right)\right)^{\frac{1}{2}} \tag{3.16}
\end{equation*}
$$

We can also easily conclude that there exist the quantities $\underline{\eta}_{1}=\underline{\eta}_{1}(\bar{\theta}(t))<0$ and $\bar{\eta}_{1}=\bar{\eta}_{1}(\bar{\theta}(t))>0$, such that

$$
\begin{equation*}
\int_{\underline{\eta}_{1}}^{0} \sqrt{e^{\eta}-1-\eta} d \eta=\int_{0}^{\bar{\eta}_{1}} \sqrt{e^{\eta}-1-\eta} d \eta=K_{2}(\bar{\theta}(t)) \tag{3.17}
\end{equation*}
$$

where $K_{2}(\bar{\theta}(t))$ is defined by (3.16). Comparing (3.15) and (3.17) we conclude that

$$
\begin{equation*}
\underline{u}_{1}=\exp \underline{\eta}_{1} \leq \rho^{-1}(x, \tau) \leq \exp \bar{\eta}_{1}=\bar{u}_{1} \tag{3.18}
\end{equation*}
$$

for $(x, \tau) \in \mathbb{R} \times] 0, t[$.
Now we find some estimates for the derivatives of the functions $v$ and $\theta$.

Lemma 3.3. For each $t>0$ we have

$$
\begin{align*}
& \frac{1}{2} \int_{\mathbb{R}}\left(\left(\frac{\partial v}{\partial x}\right)^{2}+\left(\frac{\partial \theta}{\partial x}\right)^{2}\right) d x \\
& \quad+\frac{8 \bar{u}_{1}}{\left(D \underline{u}_{1}\right)^{2}} \int_{0}^{t}\left(\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x\right)^{2}\left(K_{4}(\bar{\theta}(\tau))-\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x\right) d \tau  \tag{3.19}\\
& \quad+\frac{D}{8} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} \theta}{\partial x^{2}}\right)^{2} d x d \tau \leq K_{3}(\bar{\theta}(t))
\end{align*}
$$

where

$$
\begin{gather*}
K_{4}(\bar{\theta}(t))=\left(\frac{D \underline{u}_{1}}{16 \bar{u}_{1} \sqrt{E_{1}}}\right)^{2}  \tag{3.20}\\
K_{3}(\bar{\theta}(t))=E_{1} \frac{\bar{u}_{1}}{\underline{u}_{1}}(1+\bar{\theta}(t)) \bar{\theta}(t)\left(\left(16 K_{1}(\bar{\theta}(t))\right)^{2} \frac{\bar{u}_{1}}{\underline{u}_{1}}+\frac{K^{2}}{D} \bar{\theta}(t)\right)  \tag{3.21}\\
+2 K \bar{\theta}(t) K_{1}(\bar{\theta}(t))+E_{2}
\end{gather*}
$$

Proof. Multiplying equations (2.2) and (2.3), respectively, by $-\frac{\partial^{2} v}{\partial x^{2}}$ and $-\frac{\partial^{2} \theta}{\partial x^{2}}$, integrating over $\left.\mathbb{R} \times\right] 0, t[$ and using the following equality

$$
\begin{equation*}
-\int_{0}^{t} \int_{\mathbb{R}} \frac{\partial v}{\partial t} \frac{\partial^{2} v}{\partial x^{2}} d x d \tau=\left.\frac{1}{2} \int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x\right|_{0} ^{t} \tag{3.22}
\end{equation*}
$$

that is satisfied for the function $\theta$ as well, after addition of the obtained equalities we find that

$$
\begin{aligned}
\frac{1}{2} & \left.\int_{\mathbb{R}}\left(\left(\frac{\partial v}{\partial x}\right)^{2}+\left(\frac{\partial \theta}{\partial x}\right)^{2}\right) d x\right|_{0} ^{t}+\int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x d \tau \\
& +D \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} \theta}{\partial x^{2}}\right)^{2} d x d \tau \\
= & -\int_{0}^{t} \int_{\mathbb{R}} \frac{\partial \rho}{\partial x} \frac{\partial v}{\partial x} \frac{\partial^{2} v}{\partial x^{2}} d x d \tau+K \int_{0}^{t} \int_{\mathbb{R}} \rho \frac{\partial \theta}{\partial x} \frac{\partial^{2} v}{\partial x^{2}} d x d \tau \\
& +K \int_{0}^{t} \int_{\mathbb{R}} \theta \frac{\partial \rho}{\partial x} \frac{\partial^{2} v}{\partial x^{2}} d x d \tau+K \int_{0}^{t} \int_{\mathbb{R}} \rho \theta \frac{\partial v}{\partial x} \frac{\partial^{2} \theta}{\partial x^{2}} d x d \tau \\
& -\int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial v}{\partial x}\right)^{2} \frac{\partial^{2} \theta}{\partial x^{2}} d x d \tau-D \int_{0}^{t} \int_{\mathbb{R}} \frac{\partial \rho}{\partial x} \frac{\partial \theta}{\partial x} \frac{\partial^{2} \theta}{\partial x^{2}} d x d \tau
\end{aligned}
$$

Using (3.18), (3.11) and the following inequality

$$
\begin{equation*}
\left(\frac{\partial v}{\partial x}\right)^{2} \leq 2\left(\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2}\right)^{\frac{1}{2}}\left(\int_{\mathbb{R}}\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2}\right)^{\frac{1}{2}} \tag{3.24}
\end{equation*}
$$

that holds for the function $\frac{\partial \theta}{\partial x}$ as well, and applying the Young's inequality with a sufficiently small parameter on the right-hand side of (3.23) we come to the estimates as follows:
(3.25)

$$
\begin{align*}
& \left|\int_{0}^{t} \int_{\mathbb{R}} \frac{\partial \rho}{\partial x} \frac{\partial v}{\partial x} \frac{\partial^{2} v}{\partial x^{2}} d x d \tau\right| \\
& \leq \int_{0}^{t} \int_{\mathbb{R}} \frac{1}{\rho}\left(\frac{\partial \rho}{\partial x}\right)^{2}\left(\frac{\partial v}{\partial x}\right)^{2} d x d \tau+\frac{1}{4} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x d \tau \\
& \leq \frac{2 \bar{u}_{1}^{\frac{1}{2}}}{\underline{u}_{1}} \int_{0}^{t}\left(\int_{\mathbb{R}} \frac{1}{\rho^{2}}\left(\frac{\partial \rho}{\partial x}\right)^{2} d x\right)\left(\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x\right)^{\frac{1}{2}}\left(\int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x\right)^{\frac{1}{2}} d \tau \\
& \quad+\frac{1}{4} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x d \tau \\
& \leq \frac{5}{16} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x d \tau+\left(16 K_{1}(\bar{\theta}(t))\right)^{2} \frac{\bar{u}_{1}}{u_{1}^{2}} \int_{0}^{t} \int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x d \tau \\
& \leq \frac{5}{16} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x d \tau \\
& \quad+\left(16 K_{1}(\bar{\theta}(t)) \frac{\bar{u}_{1}}{\underline{u}_{1}}\right)^{2} \bar{\theta}(t) \int_{0}^{t} \int_{\mathbb{R}} \frac{\rho}{\theta}\left(\frac{\partial v}{\partial x}\right)^{2} d x d \tau, \\
& (3.26) \\
& \quad \leq\left.\int_{0}^{t} \int_{\mathbb{R}} \rho \frac{\partial \theta}{\partial x} \frac{\partial^{2} v}{\partial x^{2}} d x d \tau\right|^{2} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial \theta}{\partial x}\right)^{2} d x d \tau+\frac{1}{4} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x d \tau  \tag{3.26}\\
& \quad \leq \frac{K^{2} \bar{u}_{1} \bar{\theta}^{2}(t)}{\underline{u}_{1}} \int_{0}^{t} \int_{\mathbb{R}} \frac{\rho}{\theta^{2}}\left(\frac{\partial \theta}{\partial x}\right)^{2} d x d \tau+\frac{1}{4} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x d \tau \\
& \quad \\
& \quad \leq K^{2} \bar{\theta}(t) \int_{0}^{t} \int_{\mathbb{R}} \theta \rho^{3}\left(\frac{\partial}{\partial x}\left(\frac{1}{\rho}\right)\right)^{2} d x d \tau+\frac{1}{4} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x d \tau  \tag{3.27}\\
& (3.27) \\
& \left.\quad \leq \int_{0}^{t} \int_{\mathbb{R}} \theta \frac{\partial \rho}{\partial x} \frac{\partial^{2} v}{\partial x^{2}} d x d \tau \right\rvert\, \\
& \quad
\end{align*}
$$

$$
\begin{align*}
& \left|K \int_{0}^{t} \int_{\mathbb{R}} \rho \theta \frac{\partial v}{\partial x} \frac{\partial^{2} \theta}{\partial x^{2}} d x d \tau\right| \\
& \quad \leq \frac{K^{2}}{D} \int_{0}^{t} \int_{\mathbb{R}} \theta^{2} \rho\left(\frac{\partial v}{\partial x}\right)^{2} d x d \tau+\frac{D}{4} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} \theta}{\partial x^{2}}\right)^{2} d x d \tau  \tag{3.28}\\
& \quad \leq \frac{K^{2} \bar{u}_{1} \bar{\theta}^{3}(t)}{D \underline{u}_{1}} \int_{0}^{t} \int_{\mathbb{R}} \frac{\rho}{\theta}\left(\frac{\partial v}{\partial x}\right)^{2} d x d \tau+\frac{D}{4} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} \theta}{\partial x^{2}}\right)^{2} d x d \tau,
\end{align*}
$$

$$
\left|\int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial v}{\partial x}\right)^{2} \frac{\partial^{2} \theta}{\partial x^{2}} d x d \tau\right|
$$

$$
\leq \frac{1}{D} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial v}{\partial x}\right)^{4} d x d \tau+\frac{D}{4} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} \theta}{\partial x^{2}}\right)^{2} d x d \tau
$$

$$
\begin{equation*}
\leq \frac{2 \bar{u}_{1}^{\frac{1}{2}}}{D \underline{u}_{1}} \int_{0}^{t}\left(\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x\right)^{\frac{3}{2}}\left(\int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x\right)^{\frac{1}{2}} d \tau \tag{3.29}
\end{equation*}
$$

$$
+\frac{D}{4} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} \theta}{\partial x^{2}}\right)^{2} d x d \tau
$$

$$
\leq \frac{8 \bar{u}_{1}}{D^{2} \underline{u}_{1}^{2}} \int_{0}^{t}\left(\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x\right)^{3} d \tau+\frac{1}{8} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x d \tau
$$

$$
+\frac{D}{4} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} \theta}{\partial x^{2}}\right)^{2} d x d \tau
$$

(3.30)

$$
\begin{aligned}
& \left|D \int_{0}^{t} \int_{\mathbb{R}} \frac{\partial \rho}{\partial x} \frac{\partial \theta}{\partial x} \frac{\partial^{2} \theta}{\partial x^{2}} d x d \tau\right| \\
& \leq D \int_{0}^{t} \int_{\mathbb{R}} \frac{1}{\rho}\left(\frac{\partial \rho}{\partial x}\right)^{2}\left(\frac{\partial \theta}{\partial x}\right)^{2} d x d \tau+\frac{D}{4} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} \theta}{\partial x^{2}}\right)^{2} d x d \tau \\
& \leq \frac{2 D \bar{u}_{1}^{\frac{1}{2}}}{\underline{u}_{1}} \int_{0}^{t}\left(\int_{\mathbb{R}} \rho\left(\frac{\partial^{2} \theta}{\partial x^{2}}\right)^{2} d x\right)^{\frac{1}{2}}\left(\int_{\mathbb{R}}\left(\frac{\partial \theta}{\partial x}\right)^{2} d x\right)^{\frac{1}{2}} \int_{\mathbb{R}} \frac{1}{\rho^{2}}\left(\frac{\partial \rho}{\partial x}\right)^{2} d x d \tau \\
& \quad+\frac{D}{4} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} \theta}{\partial x^{2}}\right)^{2} d x d \tau \\
& \leq \frac{128 D K_{1}^{2}(\bar{\theta}(t)) \bar{\theta}^{2}(t) \bar{u}_{1}^{2}}{u_{1}^{2}} \int_{0}^{t} \int_{\mathbb{R}} \frac{\rho}{\theta^{2}}\left(\frac{\partial \theta}{\partial x}\right)^{2} d x d \tau \\
& \quad+\frac{3 D}{8} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} \theta}{\partial x^{2}}\right)^{2} d x d \tau .
\end{aligned}
$$

Taking into account (3.25)-(3.30), (3.1), (3.11) and (2.19) from (3.23) follows

$$
\begin{align*}
\frac{1}{2} \int_{\mathbb{R}} & {\left[\left(\frac{\partial v}{\partial x}\right)^{2}+\left(\frac{\partial \theta}{\partial x}\right)^{2}\right] d x+\frac{1}{16} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x d \tau } \\
& +\frac{D}{8} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} \theta}{\partial x^{2}}\right)^{2} d x d \tau  \tag{3.31}\\
\leq & \frac{8 \bar{u}_{1}}{D^{2} \underline{u}_{1}^{2}} \int_{0}^{t}\left(\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x\right)^{3} d \tau+K_{3}(\bar{\theta}(t)),
\end{align*}
$$

where $K_{3}(\bar{\theta}(t))$ is defined by (3.21). We also get another important inequality by estimating the integral $\int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x$. We have

$$
\begin{align*}
\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x & =-\int_{\mathbb{R}} v \frac{\partial^{2} v}{\partial x^{2}} d x \\
& \leq \bar{u}_{1}^{\frac{1}{2}}\left(\int_{\mathbb{R}} v^{2} d x\right)^{\frac{1}{2}}\left(\int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x\right)^{\frac{1}{2}}  \tag{3.32}\\
& \leq\left(2 E_{1} \bar{u}_{1}\right)^{\frac{1}{2}}\left(\int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x\right)^{\frac{1}{2}}
\end{align*}
$$

Consequently,

$$
\begin{equation*}
\int_{\mathbb{R}} \rho\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x \geq\left(2 E_{1} \bar{u}_{1}\right)^{-1}\left(\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x\right)^{2} \tag{3.33}
\end{equation*}
$$

Inserting (3.33) into (3.31) we obtain

$$
\begin{aligned}
& \frac{1}{2} \int_{\mathbb{R}}\left[\left(\frac{\partial v}{\partial x}\right)^{2}+\left(\frac{\partial \theta}{\partial x}\right)^{2}\right] d x \\
& \quad+\frac{8 \bar{u}_{1}}{D^{2} \underline{u}_{1}^{2}} \int_{0}^{t}\left(\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x\right)^{2}\left[\left(\frac{D \underline{u}_{1}}{16 \bar{u}_{1} E_{1}^{\frac{1}{2}}}\right)^{2}-\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x\right] d \tau \\
& \quad+\frac{D}{8} \int_{0}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} \theta}{\partial x^{2}}\right)^{2} d x d \tau \leq K_{3}(\bar{\theta}(t))
\end{aligned}
$$

and (3.18) is satisfied.
Similarly as in [4], in the continuation we use the above results, as well as the conditions of Theorem 2.1. We derive the estimates for the solution $(\rho, v, \theta)$ of problem (2.1)-(2.8) defined by (2.9)-(2.12) in the domain $\Pi=$ $\mathbb{R} \times] 0, T[$, for arbitrary $T>0$.

Taking into account assumption (2.21) and the fact that $\theta \in C(\bar{\Pi})$ (see (2.17)) we have the following alternatives: either

$$
\begin{equation*}
\sup _{(x, t) \in \Pi} \theta(x, t)=\bar{\theta}(T) \leq M_{1} \tag{3.35}
\end{equation*}
$$

or there exists $t_{1}, 0<t_{1}<T$, such that

$$
\begin{equation*}
\bar{\theta}(t)<M_{1} \text { for } 0 \leq t<t_{1}, \quad \bar{\theta}\left(t_{1}\right)=M_{1} . \tag{3.36}
\end{equation*}
$$

Now we assume that (3.36) is satisfied and we will show later, that because of the choice of the constants $E_{1}, E_{2}, E_{3}$ and $M_{1}$ (the conditions of Theorem 2.1), (3.36) is impossible.

Because $K_{2}(\bar{\theta}(t))$, defined by (3.16), increases with increasing $\bar{\theta}(t)$ we can easily conclude that

$$
\begin{equation*}
K_{2}(\bar{\theta}(t))<K_{2}\left(M_{1}\right) \text { for } 0 \leq t<t_{1} \tag{3.37}
\end{equation*}
$$

and $K_{2}\left(M_{1}\right)=E_{5}$. Therefore we have

$$
\begin{equation*}
\underline{u}<\underline{u}_{1}(\bar{\theta}(t)), \quad \bar{u}>\bar{u}_{1}(\bar{\theta}(t)) \tag{3.38}
\end{equation*}
$$

where $\underline{u}, \bar{u}$ and $\underline{u}_{1}(\bar{\theta}(t)), \bar{u}_{1}(\bar{\theta}(t))$ are defined by (2.22)-(2.23) and (3.17)(3.18), respectively. The quantity $K_{4}(\bar{\theta}(t))$, defined by (3.20), decreases with increasing $\bar{\theta}(t)$ and for $\bar{\theta}\left(t_{1}\right)=M_{1}$ it becomes

$$
\begin{equation*}
K_{4}\left(M_{1}\right)=\left(\frac{D \underline{u}}{16 \bar{u} E_{1}^{\frac{1}{2}}}\right)^{2} \tag{3.39}
\end{equation*}
$$

Taking into account the assumption (2.25) of Theorem 2.1 and the following inclusion (See (2.16))

$$
\begin{equation*}
\frac{\partial v}{\partial x} \in \mathrm{C}\left([0, T] ; \mathrm{L}^{2}(\mathbb{R})\right) \tag{3.40}
\end{equation*}
$$

we have again two alternatives: either

$$
\begin{equation*}
\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2}(x, t) d x \leq K_{4}\left(M_{1}\right) \text { for } t \in\left[0, t_{1}\right] \tag{3.41}
\end{equation*}
$$

or there exists $t_{2}, \quad 0<t_{2}<t_{1}$, such that

$$
\begin{equation*}
\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2}(x, t) d x<K_{4}\left(M_{1}\right) \text { for } 0 \leq t<t_{2} \tag{3.42}
\end{equation*}
$$

and

$$
\begin{equation*}
\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2}\left(x, t_{2}\right) d x=K_{4}\left(M_{1}\right) \text { for } t_{2}<t_{1} \tag{3.43}
\end{equation*}
$$

Now, we assume that (3.42)-(3.43) are satisfied. Then we have

$$
\begin{equation*}
\bar{\theta}(t)<M_{1}, \quad K_{4}\left(M_{1}\right)<K_{4}(\bar{\theta}(t)) \text { for } t \in\left[0, t_{2}\right] . \tag{3.44}
\end{equation*}
$$

Taking into account (3.44), from (3.19), for $t=t_{2}$, we obtain

$$
\begin{equation*}
\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2}(x, t) d x \leq 2 K_{3}\left(\bar{\theta}\left(t_{2}\right)\right), \quad 0 \leq t \leq t_{2} \tag{3.45}
\end{equation*}
$$

Since $K_{3}(\bar{\theta}(t))$, defined by (3.21), increases with the increase of $\bar{\theta}(t)$, it holds

$$
\begin{equation*}
K_{3}(\bar{\theta}(t)) \leq K_{3}\left(M_{1}\right), \quad t \in\left[0, t_{2}\right] . \tag{3.46}
\end{equation*}
$$

Using condition (2.26) we get

$$
\begin{equation*}
2 K_{3}(\bar{\theta}(t))<K_{4}\left(M_{1}\right), \quad t \in\left[0, t_{2}\right], \tag{3.47}
\end{equation*}
$$

and conclude that

$$
\begin{equation*}
\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2}\left(x, t_{2}\right) d x<K_{4}\left(M_{1}\right) \tag{3.48}
\end{equation*}
$$

This inequality contradicts (3.43). Consequently, the only case possible is when

$$
\begin{equation*}
t_{2}=t_{1} \tag{3.49}
\end{equation*}
$$

and then (3.41) is satisfied.
Using (3.36) and (3.41) from (3.19) we can easily obtain that

$$
\begin{equation*}
\int_{\mathbb{R}}\left(\frac{\partial \theta}{\partial x}\right)^{2}(x, t) d x<2 K_{3}\left(M_{1}\right) \text { for } 0<t \leq t_{1} \tag{3.50}
\end{equation*}
$$

Now, we introduce the function $\Psi$ by

$$
\begin{equation*}
\Psi(\theta(x, t))=\int_{1}^{\theta(x, t)} \sqrt{s-1-\ln s} d s \tag{3.51}
\end{equation*}
$$

From (2.17) follows that $\theta(x, t) \rightarrow 1$ as $|x| \rightarrow \infty$ and hence

$$
\begin{equation*}
\Psi(\theta(x, t)) \rightarrow 0 \quad \text { as } \quad|x| \rightarrow \infty \tag{3.52}
\end{equation*}
$$

Consequently,

$$
\begin{align*}
\psi(\theta(x, t)) & \leq|\psi(\theta(x, t))|=\left|\int_{1}^{\theta(x, t)} \frac{d}{d s} \psi(s) d s\right|  \tag{3.53}\\
& =\left|\int_{-\infty}^{x} \sqrt{\theta(x, t)-1-\ln \theta(x, t)} \frac{\partial \theta(x, t)}{\partial x} d x\right| \\
& \leq\left(\int_{\mathbb{R}}(\theta(x, t)-1-\ln \theta(x, t)) d x\right)^{\frac{1}{2}}\left(\int_{\mathbb{R}}\left(\frac{\partial \theta}{\partial x}\right)^{2}(x, t) d x\right)^{\frac{1}{2}}
\end{align*}
$$

Taking into account (3.36), (3.50) and (3.1) from (3.53) we get

$$
\begin{equation*}
\max _{0 \leq \theta(x, t) \leq M_{1}} \psi(\theta(x, t))=\psi\left(\bar{\theta}\left(t_{1}\right)\right)=\psi\left(M_{1}\right) \leq\left(2 K_{3}\left(M_{1}\right) E_{1}\right)^{\frac{1}{2}} \tag{3.54}
\end{equation*}
$$

or

$$
\begin{equation*}
\int_{1}^{M_{1}} \sqrt{s-1-\ln s} d s-\left(2 K_{3}\left(M_{1}\right) E_{1}\right)^{\frac{1}{2}} \leq 0 \tag{3.55}
\end{equation*}
$$

Since this inequality contradicts (2.26), it remains to assume that $t_{1}=T$. Hence we have

Lemma 3.4. For each $T>0$ we have

$$
\begin{gather*}
\theta(x, t) \leq M_{1}, \quad(x, t) \in \Pi  \tag{3.56}\\
\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2}(x, t) d x \leq K_{4}\left(M_{1}\right), \quad 0 \leq t \leq T  \tag{3.57}\\
\int_{\mathbb{R}}\left(\frac{\partial \theta}{\partial x}\right)^{2}(x, t) d x \leq 2 K_{3}\left(M_{1}\right), \quad 0 \leq t \leq T \tag{3.58}
\end{gather*}
$$

Proof. These conclusions follow from (3.36), (3.41) and (3.50) directly.

Lemma 3.5. The following inequalities hold true:

$$
\begin{gather*}
0<\underline{u} \leq \frac{1}{\rho(x, t)} \leq \bar{u}, \quad(x, t) \in \Pi  \tag{3.59}\\
\sup _{(x, t) \in \Pi}|v(x, t)| \leq \sqrt{8 E_{1} K_{4}\left(M_{1}\right)},  \tag{3.60}\\
\theta(x, t) \geq h>0, \quad(x, t) \in \Pi \tag{3.61}
\end{gather*}
$$

where $\underline{u}$ and $\bar{u}$ are defined by (2.22)-(2.24) and a constant $h$ depends only on the data of problem (2.1)-(2.8).

Proof. Because the quantity $\underline{u}_{1}(\bar{\theta}(t))$ in Lemma 3.2 decreases with increasing $\bar{\theta}(t)$, while $\bar{u}_{1}(\bar{\theta}(t))$ increases, it follows from (3.2) and (3.56) that (3.59) is satisfied.

Using the inequality

$$
\begin{equation*}
v^{2}=2 \int_{-\infty}^{x} v \frac{\partial v}{\partial x} d x \leq 2\left(\int_{\mathbb{R}} v^{2} d x\right)^{\frac{1}{2}}\left(\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x\right)^{\frac{1}{2}} \tag{3.62}
\end{equation*}
$$

and estimations (3.1) and (3.57) we get immediately (3.60). From (3.53), $(3.56),(3.58)$ and (3.1) we have for $\theta(x, t) \leq 1$ that the following holds

$$
\begin{equation*}
\int_{\theta(x, t)}^{1} \sqrt{s-1-\ln s} d s \leq\left(2 K_{3}\left(M_{1}\right) E_{1}\right)^{\frac{1}{2}}<\int_{0}^{1} \sqrt{s-1-\ln s} d s \tag{3.63}
\end{equation*}
$$

because of (2.26). Hence we conclude that there exists the constant $h>0$ such that $\theta(x, t) \geq h$.

Lemma 3.6. For each $T>0$ we have

$$
\begin{equation*}
\int_{\mathbb{R}}\left(\frac{\partial \rho}{\partial x}\right)^{2} d x \leq K_{8}, \quad t \in[0, T] \tag{3.67}
\end{equation*}
$$

$$
\begin{equation*}
\int_{0}^{T} \int_{\mathbb{R}}\left(\frac{\partial^{2} \theta}{\partial x^{2}}\right)^{2} d x d \tau \leq K_{9} \tag{3.68}
\end{equation*}
$$

$$
\begin{align*}
& \int_{0}^{T} \int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x d \tau \leq K_{5}  \tag{3.64}\\
& \int_{0}^{T} \int_{\mathbb{R}}\left(\frac{\partial \theta}{\partial x}\right)^{2} d x d \tau \leq K_{6}  \tag{3.65}\\
& \int_{0}^{T} \int_{\mathbb{R}}\left(\frac{\partial \rho}{\partial x}\right)^{2} d x d \tau \leq K_{7} \tag{3.66}
\end{align*}
$$

$$
\begin{equation*}
\int_{0}^{T} \int_{\mathbb{R}}\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x d \tau \leq K_{10} \tag{3.69}
\end{equation*}
$$

where the constants $K_{5}, K_{6}, K_{7}, K_{8}, K_{9}, K_{10} \in \mathbb{R}^{+}$are independent of $T$.
Proof. Taking into account (3.59) and (3.61) from (3.1), (3.11), (3.19) and (3.31) we get easily (3.64)-(3.69).

## 4. Proof of Theorem 2.1

In the following we use the results of Section 3. It is important to remark that all the estimates obtained above are preserved in the domain $\Pi=\mathbb{R} \times] 0, T[$ for each $T>0$.

The conclusions of Theorem 2.1 are immediate consequences of the following lemmas.

Lemma 4.1. The following relations hold true:

$$
\begin{equation*}
\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2}(x, t) d x \rightarrow 0, \quad \int_{\mathbb{R}}\left(\frac{\partial \theta}{\partial x}\right)^{2}(x, t) d x \rightarrow 0 \tag{4.1}
\end{equation*}
$$

when $t \rightarrow \infty$.
Proof. Let $\varepsilon>0$ be arbitrary. With the help of (3.64), (3.65) and (3.66) we conclude that there exists $t_{0}>0$ such that

$$
\begin{equation*}
\int_{t_{0}}^{t} \int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x d \tau<\varepsilon, \int_{t_{0}}^{t} \int_{\mathbb{R}}\left(\frac{\partial \theta}{\partial x}\right)^{2} d x d \tau<\varepsilon, \int_{t_{0}}^{t} \int_{\mathbb{R}}\left(\frac{\partial \rho}{\partial x}\right)^{2} d x d \tau<\varepsilon \tag{4.2}
\end{equation*}
$$

for $t>t_{0}$, and

$$
\begin{equation*}
\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2}\left(x, t_{0}\right) d x<\varepsilon, \quad \int_{\mathbb{R}}\left(\frac{\partial \theta}{\partial x}\right)^{2}\left(x, t_{0}\right) d x<\varepsilon \tag{4.3}
\end{equation*}
$$

Similarly to (3.19), we have

$$
\begin{aligned}
\frac{1}{2} \int_{\mathbb{R}} & {\left[\left(\frac{\partial v}{\partial x}\right)^{2}+\left(\frac{\partial \theta}{\partial x}\right)^{2}\right] d x } \\
& +\frac{8 \bar{u}}{(D \underline{u})^{2}} \int_{t_{0}}^{t}\left(\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x\right)^{2}\left[K_{4}(\bar{\theta}(\tau))-\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x\right] d \tau \\
& +\frac{D}{8} \int_{t_{0}}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial^{2} \theta}{\partial x^{2}}\right)^{2} d x d \tau \\
\leq & \frac{1}{2} \int_{\mathbb{R}}\left[\left(\frac{\partial v}{\partial x}\right)^{2}+\left(\frac{\partial \theta}{\partial x}\right)^{2}\right]\left(x, t_{0}\right) d x+K^{2} \int_{t_{0}}^{t} \int_{\mathbb{R}} \rho\left(\frac{\partial \theta}{\partial x}\right)^{2} d x d \tau \\
& +\left(16 K_{1}(\bar{\theta}(t))\right)^{2} \frac{\bar{u}}{u^{2}} \int_{t_{0}}^{t} \int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2} d x d \tau \\
& +K^{2} \int_{t_{0}}^{t} \int_{\mathbb{R}} \frac{\theta^{2}}{\rho}\left(\frac{\partial \rho}{\partial x}\right)^{2} d x d \tau+\frac{K^{2}}{D} \int_{t_{0}}^{t} \int_{\mathbb{R}} \theta^{2} \rho\left(\frac{\partial v}{\partial x}\right)^{2} d x d \tau \\
& +\frac{128 D K_{1}^{2}(\bar{\theta}(t)) \bar{u}}{\underline{u}^{2}} \int_{t_{0}}^{t} \int_{\mathbb{R}}\left(\frac{\partial \theta}{\partial x}\right)^{2} d x d \tau .
\end{aligned}
$$

Taking into account (3.56), (3.59) and (4.2)-(4.3) from (4.4) we obtain

$$
\begin{equation*}
\int_{\mathbb{R}}\left[\left(\frac{\partial v}{\partial x}\right)^{2}+\left(\frac{\partial \theta}{\partial x}\right)^{2}\right] d x \leq K_{11} \varepsilon \text { for } t>t_{0} \tag{4.5}
\end{equation*}
$$

where $K_{11}$ depends only on the data of our problem and does not depend on $t_{0}$. Hence relations (4.1) hold.

Lemma 4.2. We have

$$
\begin{equation*}
v(x, t) \rightarrow 0, \quad \theta(x, t) \rightarrow 1 \tag{4.6}
\end{equation*}
$$

when $t \rightarrow \infty$, uniformly with respect to all $x \in \mathbb{R}$.
Proof. We have (see (3.62) and (3.53))

$$
\begin{equation*}
v^{2}(x, t) \leq 2\left(\int_{\mathbb{R}} v^{2}(x, t) d x\right)^{\frac{1}{2}}\left(\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2}(x, t) d x\right)^{\frac{1}{2}} \tag{4.7}
\end{equation*}
$$

(4.8) $|\psi(\theta(x, t))| \leq\left(\int_{\mathbb{R}}(\theta(x, t)-1-\ln \theta(x, t)) d x\right)^{\frac{1}{2}}\left(\int_{\mathbb{R}}\left(\frac{\partial \theta}{\partial x}\right)^{2} d x\right)^{\frac{1}{2}}$.

Taking into account (3.1) from (4.7) and (4.8) we get

$$
\begin{gather*}
v^{2}(x, t) \leq 2\left(2 E_{1}\right)^{\frac{1}{2}}\left(\int_{\mathbb{R}}\left(\frac{\partial v}{\partial x}\right)^{2}(x, t) d x\right)^{\frac{1}{2}}  \tag{4.9}\\
|\psi(\theta(x, t))| \leq E_{1}^{\frac{1}{2}}\left(\int_{\mathbb{R}}\left(\frac{\partial \theta}{\partial x}\right)^{2} d x\right)^{\frac{1}{2}} \tag{4.10}
\end{gather*}
$$

Using (4.1) and property (3.52) of the function $\psi$ we can easily obtain that (4.6) holds.

Lemma 4.3. We have

$$
\begin{equation*}
\int_{\mathbb{R}}\left(\frac{\partial \rho}{\partial x}\right)^{2}(x, t) d x \rightarrow 0 \tag{4.11}
\end{equation*}
$$

when $t \rightarrow \infty$.
Proof. From (3.66) and (3.69) we conclude that for $\varepsilon>0$ exists $t_{0}>0$ such that

$$
\begin{equation*}
\int_{t_{0}}^{t} \int_{\mathbb{R}}\left(\frac{\partial \rho}{\partial x}\right)^{2} d x d \tau<\varepsilon, \int_{t_{0}}^{t} \int_{\mathbb{R}}\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x d \tau<\varepsilon, \int_{\mathbb{R}}\left(\frac{\partial \rho}{\partial x}\right)^{2}\left(x, t_{0}\right) d x<\varepsilon \tag{4.12}
\end{equation*}
$$

for $t>t_{0}$. Now, from (2.1) we get the equality

$$
\begin{equation*}
\frac{\partial}{\partial t}\left(\frac{\partial}{\partial x}\left(\frac{1}{\rho}\right)\right)=\frac{\partial^{2} v}{\partial x^{2}} \tag{4.13}
\end{equation*}
$$

Multiplying (4.13) by $\frac{\partial}{\partial x}\left(\frac{1}{\rho}\right)$ and integrating over $\mathbb{R}$ and $] t_{0}, t[$ we obtain

$$
\begin{equation*}
\frac{1}{2} \int_{\mathbb{R}} \frac{1}{\rho^{4}}\left(\frac{\partial \rho}{\partial x}\right)^{2} d x=-\int_{t_{0}}^{t} \int_{\mathbb{R}} \frac{1}{\rho^{2}} \frac{\partial \rho}{\partial x} \frac{\partial^{2} v}{\partial x^{2}} d x d \tau+\frac{1}{2} \int_{\mathbb{R}} \frac{1}{\rho^{4}}\left(\frac{\partial \rho}{\partial x}\right)^{2}\left(x, t_{0}\right) d x \tag{4.14}
\end{equation*}
$$

Using the Young's inequality and (3.59) from (4.14) we find out

$$
\begin{align*}
\frac{u^{4}}{2} \int_{\mathbb{R}}\left(\frac{\partial \rho}{\partial x}\right)^{2} d x \leq & \bar{u}^{2} \int_{t_{0}}^{t} \int_{\mathbb{R}}\left(\frac{\partial \rho}{\partial x}\right)^{2} d x d \tau+\bar{u}^{2} \int_{t_{0}}^{t} \int_{\mathbb{R}}\left(\frac{\partial^{2} v}{\partial x^{2}}\right)^{2} d x d \tau  \tag{4.15}\\
& +\frac{\underline{u}^{4}}{2} \int_{\mathbb{R}}\left(\frac{\partial \rho}{\partial x}\right)^{2}\left(x, t_{0}\right) d x
\end{align*}
$$

With the help of (4.12) from (4.15) we get (4.11).
Lemma 4.4. We have

$$
\begin{equation*}
\rho(x, t) \rightarrow 1 \tag{4.16}
\end{equation*}
$$

when $t \rightarrow \infty$, uniformly with respect to $x \in \mathbb{R}$.

Proof. Similarly as for the function $\psi(\theta)$, we have

$$
\begin{align*}
\psi\left(\frac{1}{\rho}\right) & =\int_{1}^{\frac{1}{\rho}} \sqrt{s-1-\ln s} d s \\
& \leq\left(\int_{\mathbb{R}}\left(\frac{1}{\rho}-1-\ln \frac{1}{\rho}\right) d x\right)^{\frac{1}{2}}\left(\int_{\mathbb{R}} \frac{1}{\rho^{2}}\left(\frac{\partial \rho}{\partial x}\right)^{2} d x\right)^{\frac{1}{2}} \tag{4.17}
\end{align*}
$$

Taking into account (3.1) from (4.17) follows

$$
\begin{equation*}
\psi\left(\frac{1}{\rho}\right) \leq\left(K E_{1}\right)^{\frac{1}{2}} \bar{u}\left(\int_{\mathbb{R}}\left(\frac{\partial \rho}{\partial x}\right)^{2} d x\right)^{\frac{1}{2}} \tag{4.18}
\end{equation*}
$$

and with the help of (4.11) we conclude (4.16).

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[^0]:    ${ }^{1}$ Derivation of the equations (2.1)-(2.3) from the Eulerian description is given in [1], pp. 31-42.

